

Electric control of non-collinear magnets



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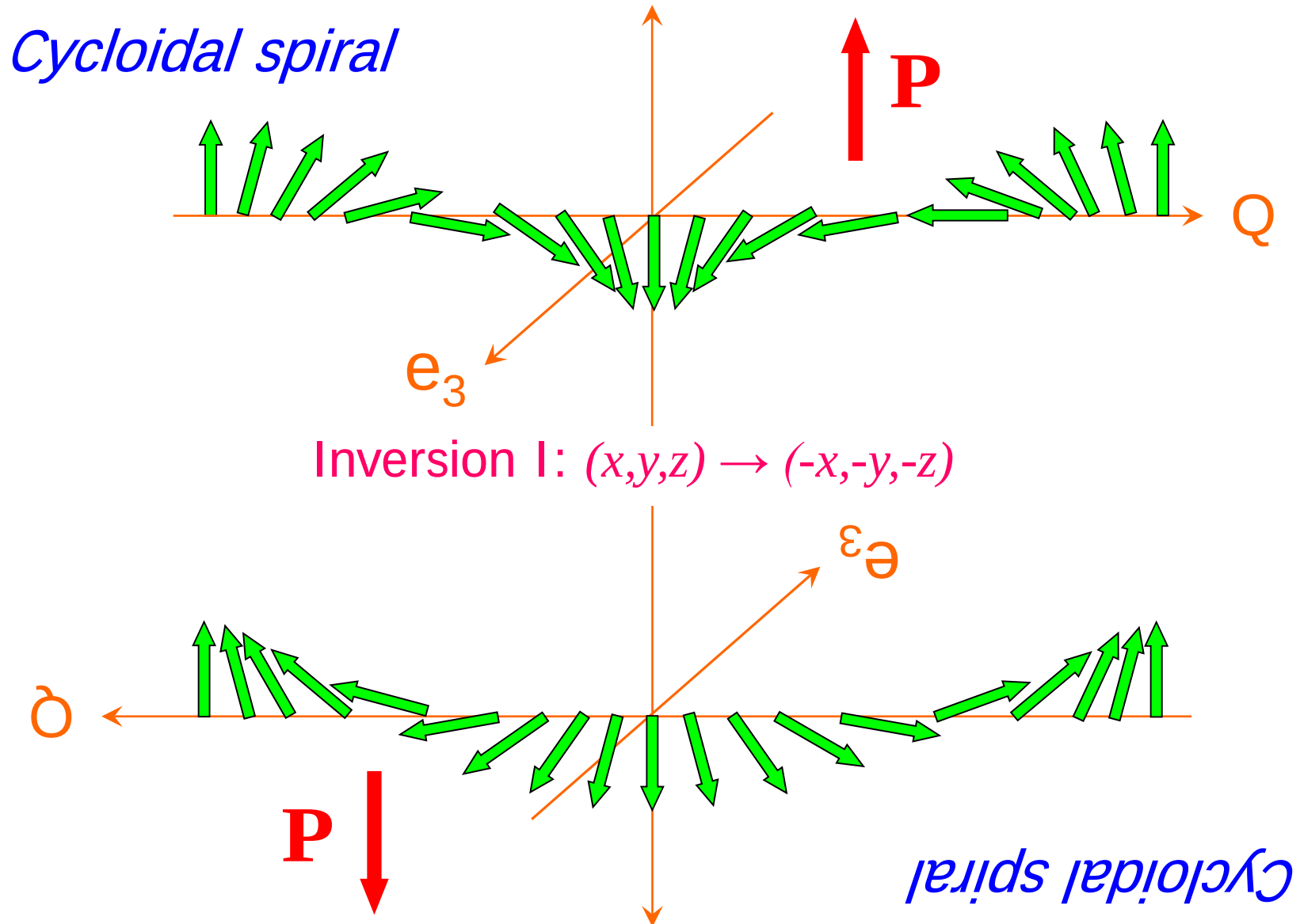
Evgenii Barts

Outline

- Symmetry of (electro-)magnons in spiral magnets
- Electric field control of polarization rotation
- Electromagnon and skyrmion mass
- Electrically excited dynamics of topological magnetic defects

Electromagnons in RMnO_3

Breaking of inversion symmetry by spin ordering



Directional dichroism

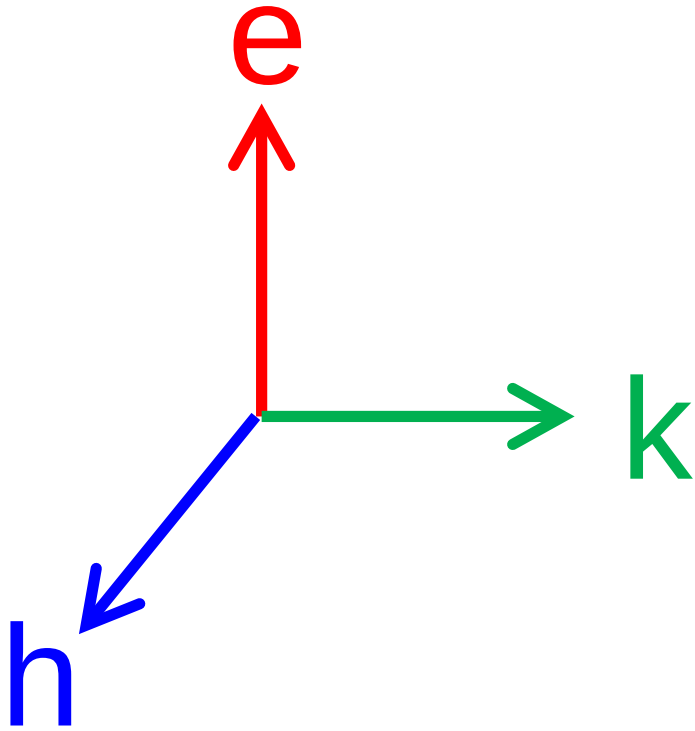
If both T and I broken

$$\delta n = \gamma \mathbf{k} \cdot [\mathbf{E} \times \mathbf{H}]$$

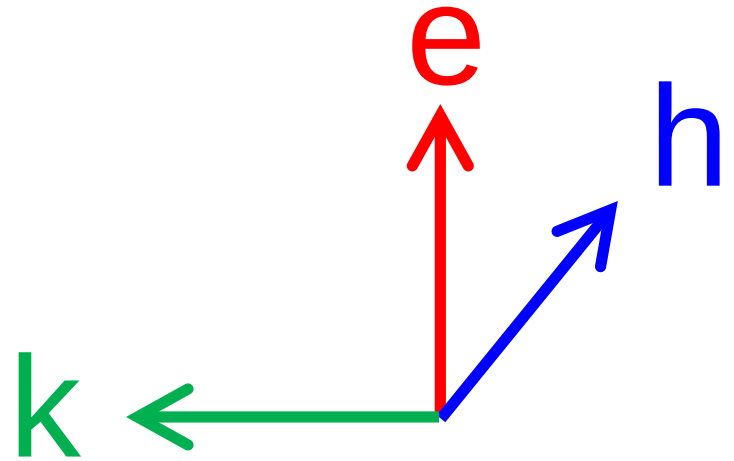
N.N. Baranova et al Sov. Phys. Usp. 20, 870 (1977)

G. L. J. A. Rikken et al Phys. Rev. Lett. 89, 133005 (2002)

Directional dichroism



$$|f_e + f_m|^2$$



$$|f_e - f_m|^2$$

Generators of Pbnm

Pbnm (# 62)

$$E \quad 1 \quad x, y, z$$

$$\tilde{m}_x \quad 2 \quad \frac{1}{2} - x, \frac{1}{2} + y, z$$

$$\tilde{m}_y \quad 3 \quad \frac{1}{2} + x, \frac{1}{2} - y, \frac{1}{2} + z$$

$$m_z \quad 4 \quad x, y, \frac{1}{2} - z$$

$$I \quad 5 \quad \bar{x}, \bar{y}, \bar{z}$$

$$I = \tilde{m}_x \tilde{m}_y m_z - (0,1,1)$$

$$\tilde{2}_x \quad 6 \quad \frac{1}{2} + x, \frac{1}{2} - y, \bar{z}$$

$$\tilde{2}_x = \tilde{m}_x I$$

$$\tilde{2}_y \quad 7 \quad \frac{1}{2} - x, \frac{1}{2} + y, \frac{1}{2} - z$$

$$\tilde{2}_y = \tilde{m}_y I$$

$$\tilde{2}_z \quad 8 \quad \bar{x}, \bar{y}, \frac{1}{2} + z$$

$$\tilde{2}_z = m_z I$$

Symmetries of the *ab* spiral

	A_x	A_y	A_z	$\cos Qy$	$\sin Qy$
$\tilde{m}_y$	+	-	+	+	-
m_z	+	+	-	+	+

$$\langle \mathbf{S} \rangle = (A_x \cos Qy, A_y \sin Qy, 0)$$

Time-reversal symmetry: $\mathbf{S} \rightarrow -\mathbf{S}$

$$T + S_y \left(\frac{\pi}{Q} \right)$$

Symmetry of (electro)magnons

	e_x	e_y	e_z	h_x	h_y	h_z
$\tilde{m}_y$	+	-	+	-	+	-
m_z	+	+	-	-	-	+

$(+, -)$ $\mathbf{e}_z$ & $\mathbf{h}_y$ $\delta\mathbf{S} = (0, 0, \delta A_z \cos Qy)$

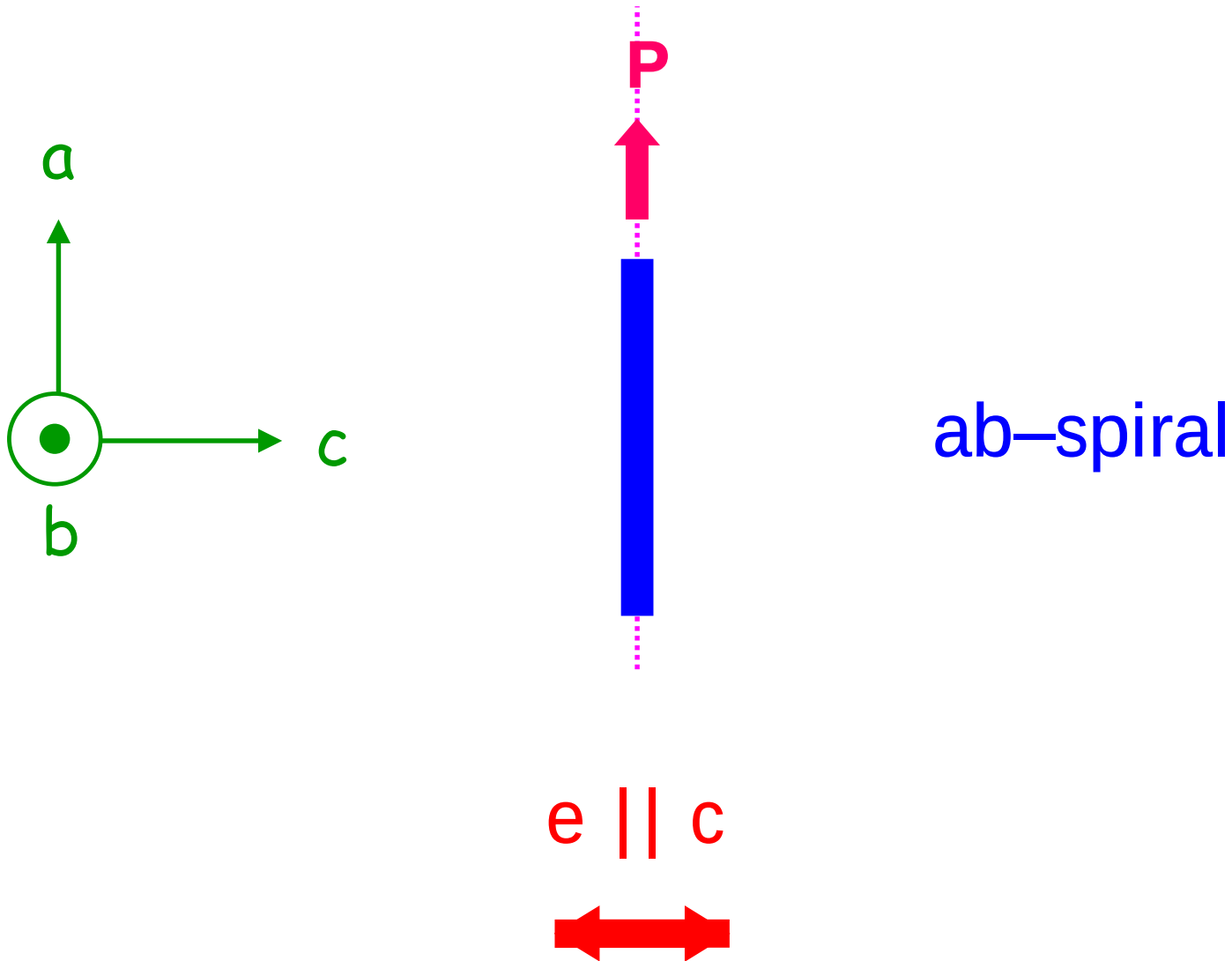
$(-, -)$ $\mathbf{h}_x$ $\delta\mathbf{S} = (0, 0, \delta A_z \sin Qy)$

$(-, +)$ $\mathbf{h}_z$ **phason**

$$\delta\mathbf{S} = (-\delta A_x \sin Qy, \delta A_y \cos Qy, 0)$$

Photo-excitation of magnons

Katsura, Balatsky & Nagaosa (2006)



Symmetry of (electro)magnons

	e_x	e_y	e_z	h_x	h_y	h_z
$\tilde{m}_y$	+	-	+	-	+	-
m_z	+	+	-	-	-	+

$(+, -)$ $\mathbf{e}_z$ & $\mathbf{h}_y$ $\delta\mathbf{S} = (0, 0, \delta A_z \cos Qy)$

$(-, -)$ $\mathbf{h}_x$ $\delta\mathbf{S} = (0, 0, \delta A_z \sin Qy)$

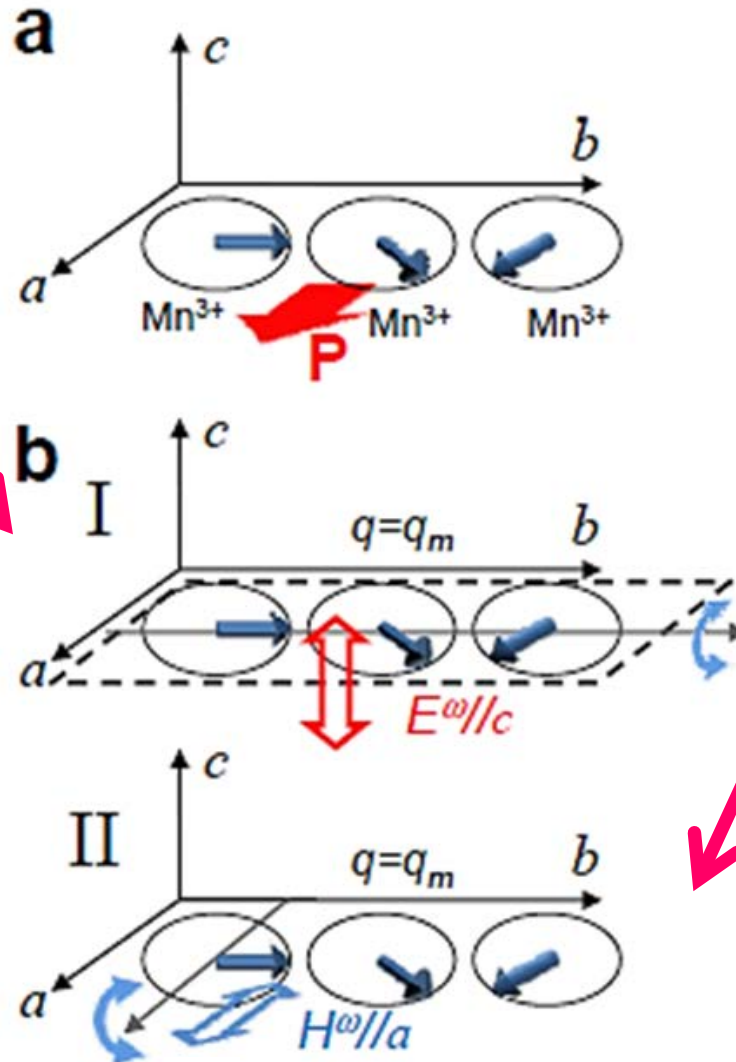
$(-, +)$ $\mathbf{h}_z$ **phason**

$$\delta\mathbf{S} = (-\delta A_x \sin Qy, \delta A_y \cos Qy, 0)$$

Collective modes in spiral state

$$e^{iQy} + e^{-iQy}$$

e-magnon
e || c



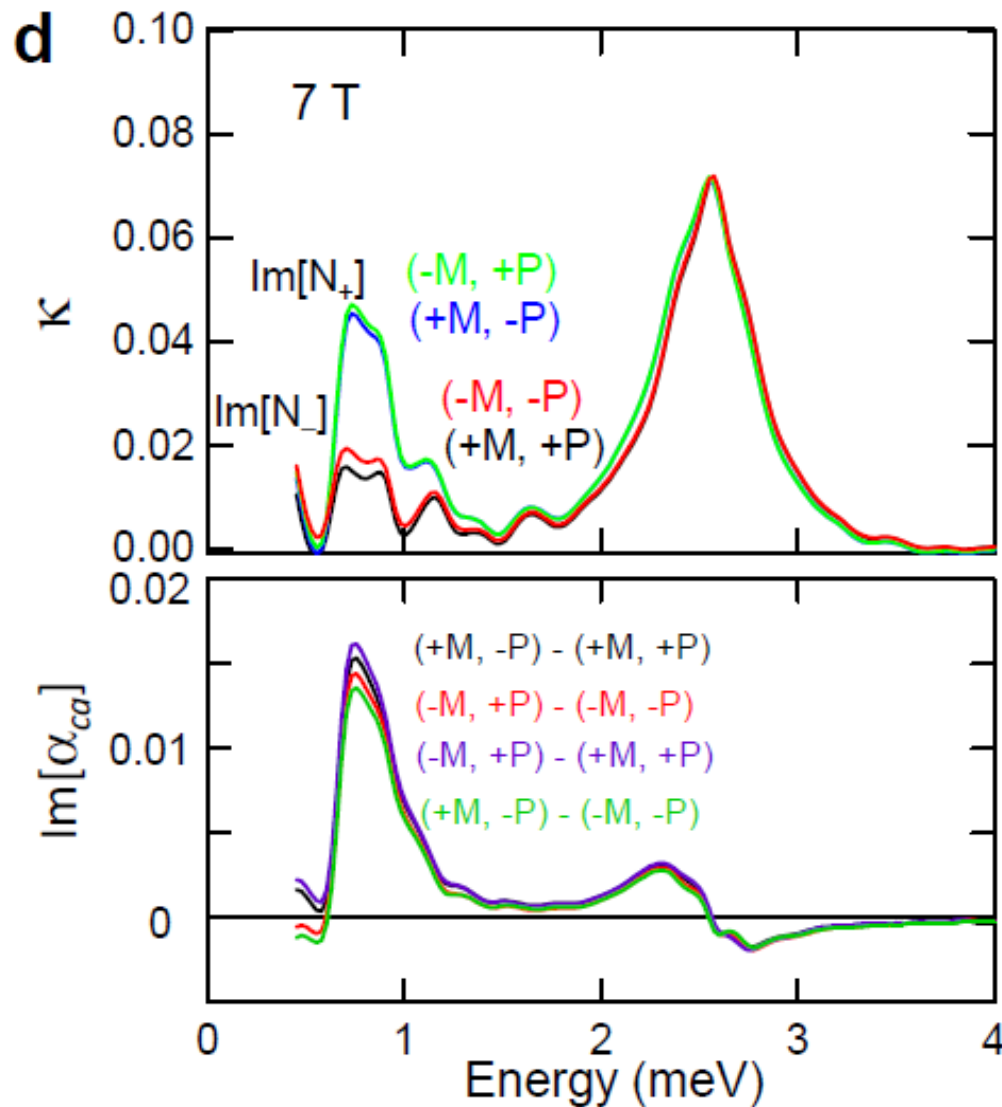
$$e^{iQy} - e^{-iQy}$$

AFMR

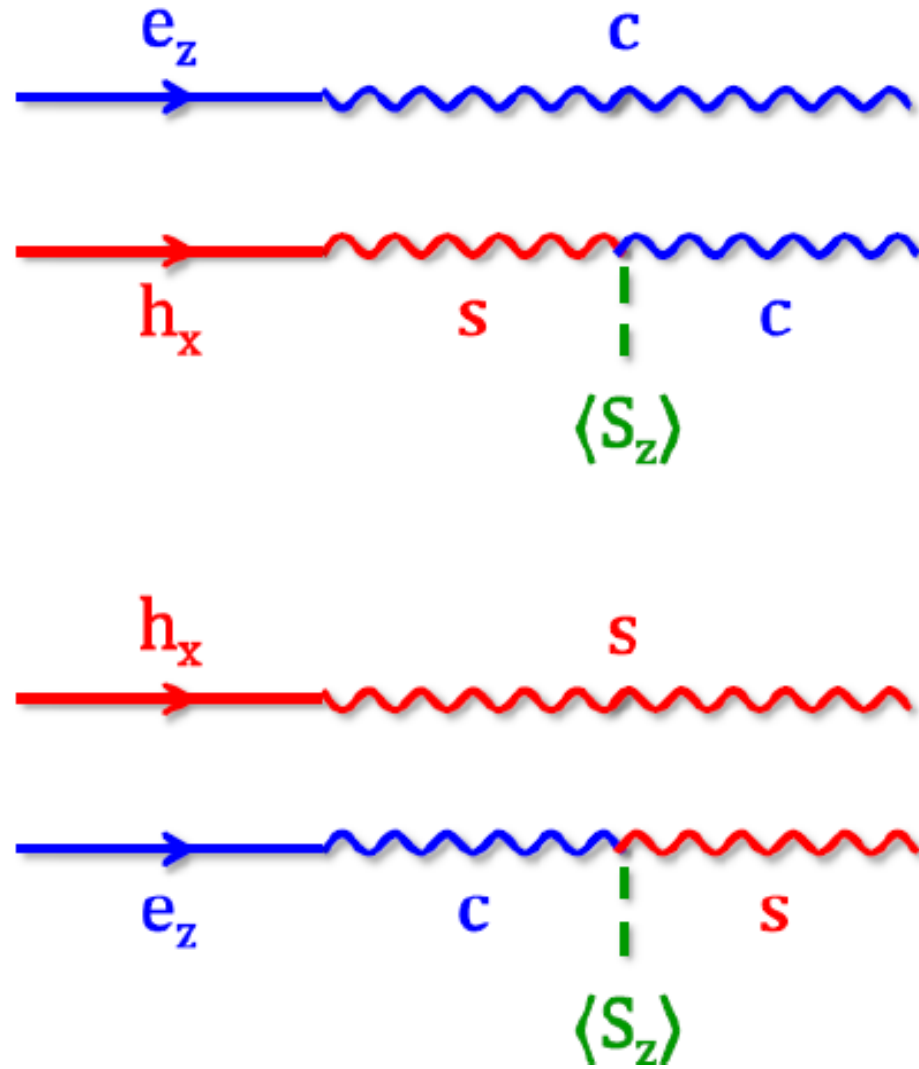
h || b

h || a

Directional dichroism in (EuY)MnO₃



Interference between e-magnon and AFM resonance



Directional dichroism and dynamical ME effect

Refraction index

$$N_{\pm} = n_{\pm} + i\kappa_{\pm}$$

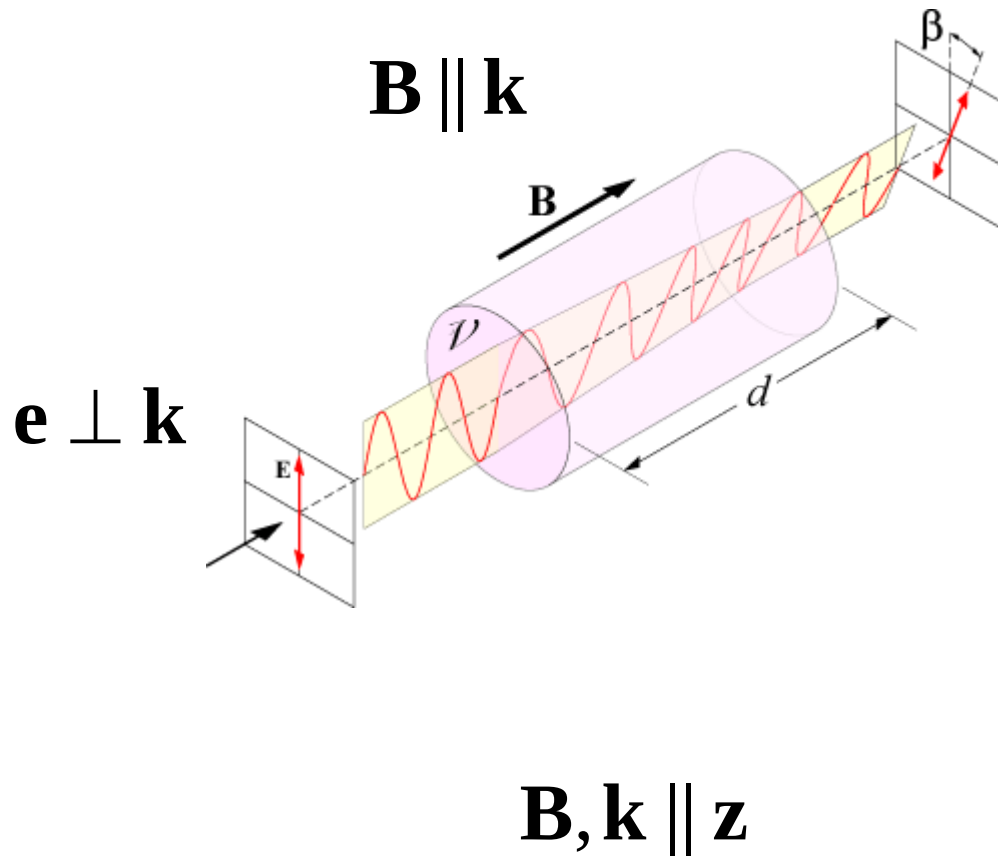
Directional dichroism

$$\kappa_{+} - \kappa_{-} = \text{Im}\left(\chi_{zx}^{em} + \chi_{xz}^{me}\right)$$

$$\delta P_z = \chi_{zx}^{em} h_x$$

$$\delta M_x = \chi_{xz}^{me} e_z$$

Faraday Effect



$$(\mathbf{k} \cdot \nabla) \mathbf{e} = C [\mathbf{B} \times \mathbf{e}]$$

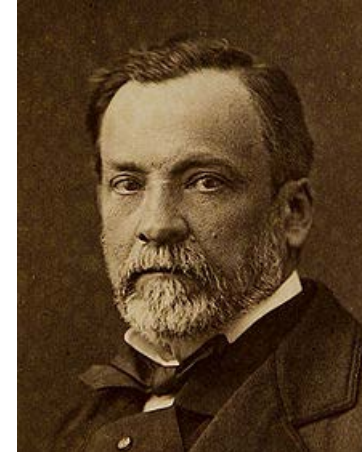
$$k_z \frac{de_x}{dz} = +CB_z e_y$$

$$k_z \frac{de_y}{dz} = -CB_z e_x$$

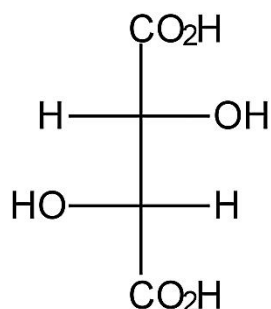


Jean Baptiste Biot
(1774 – 1862)

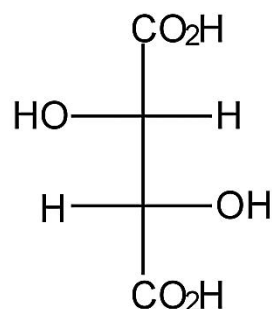
Natural Optical Activity



Louis Pasteur
(1822 – 1895)



L(+) tartaric acid



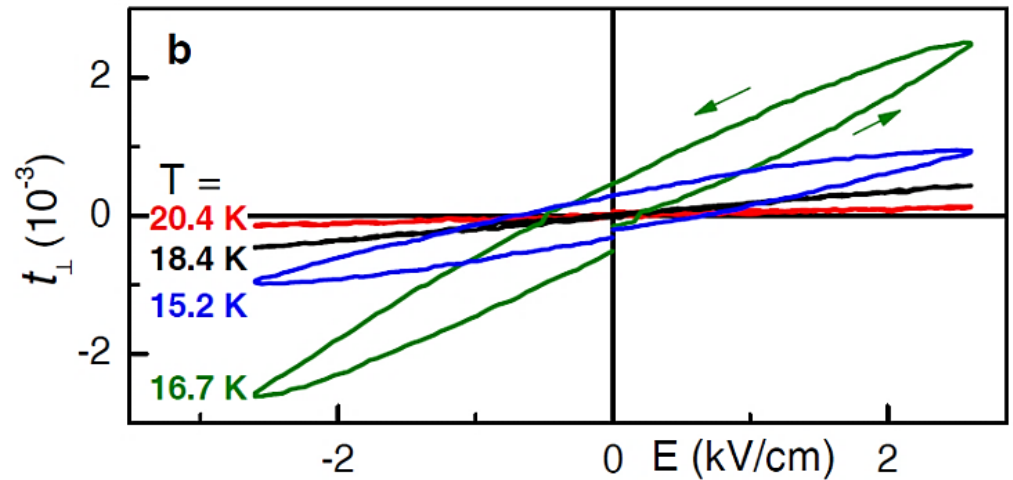
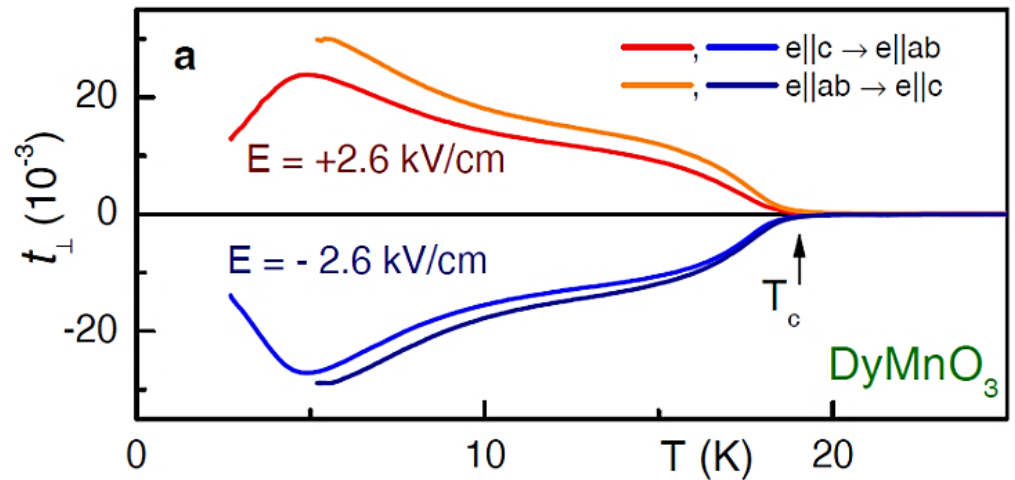
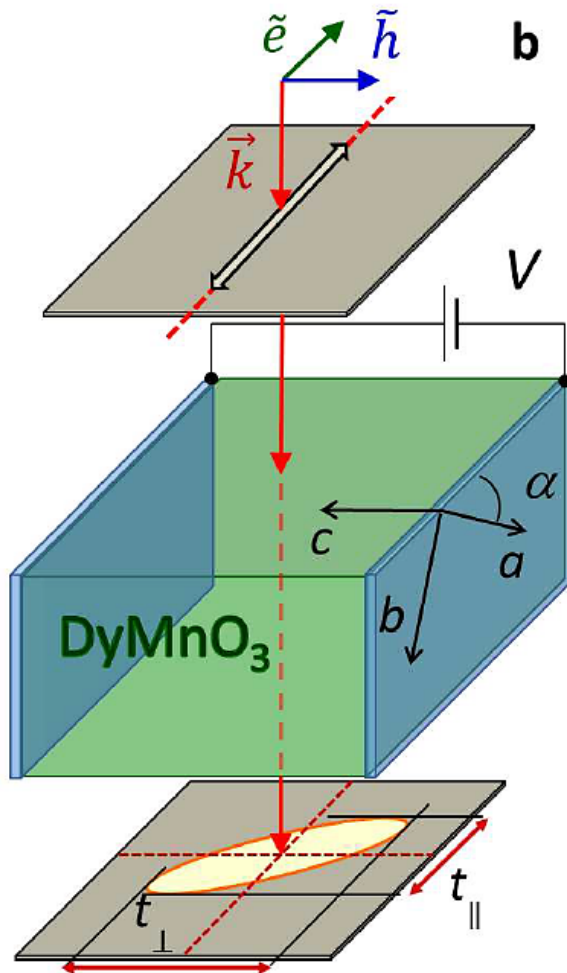
D(-) tartaric acid

$$\frac{de_x}{dz} = +Ce_y$$

$$\frac{de_y}{dz} = -Ce_x$$

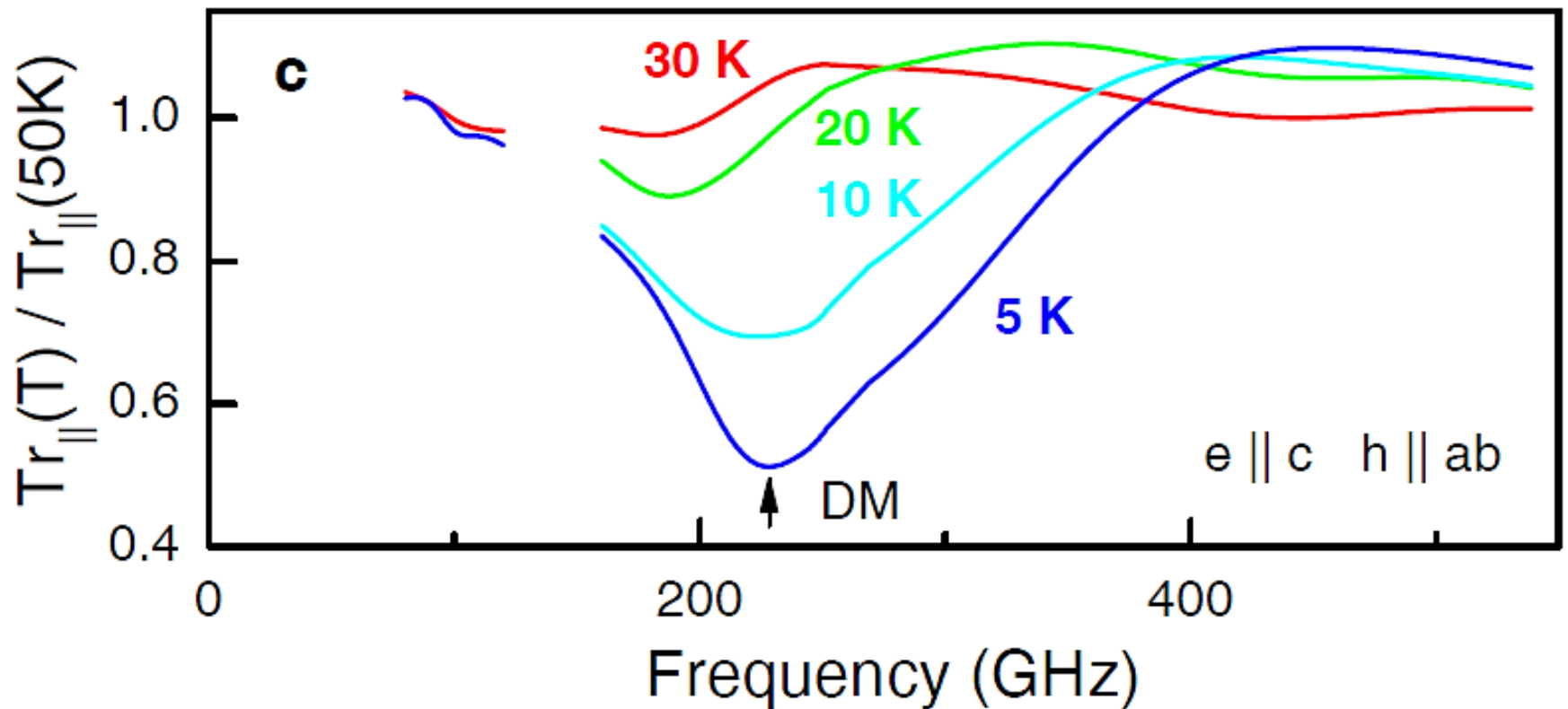
Under inversion $\mathbf{E} \rightarrow -\mathbf{E}$, $z \rightarrow -z$ equations change sign $C = -C = 0$

Electric field control of light polarization in DyMnO₃



Shuvaev et al PRL 111, 227201 (2017)

Electromagnon



Symmetries of the *bc* spiral

	A_x	A_y	A_z	$\cos Qy$	$\sin Qy$
$\tilde{m}_y$	+	-	+	+	-
$\tilde{z}_z + S_y$	-	-	+	+	-

$$\langle \mathbf{S} \rangle = (0, A_y \sin Qy, A_z \cos Qy)$$

Time-reversal symmetry: $\mathbf{S} \rightarrow -\mathbf{S}$

$$T + S_y \left(\frac{\pi}{Q} \right)$$

Symmetry of (electro)magnons

	e_x	e_y	e_z	h_x	h_y	h_z
$\tilde{m}_y$	+	-	+	-	+	-
$\tilde{z}_z + S_y$	-	-	+	-	-	+

$(+, -)$ $\mathbf{e}_x$ & $\mathbf{h}_y$ $\delta \mathbf{S} = (\delta A_x \cos Qy, 0, 0)$

$(-, +)$ $\mathbf{h}_z$ $\delta \mathbf{S} = (\delta A_x \sin Qy, 0, 0)$

$(-, -)$ $\mathbf{h}_x$ **phason**

$$\delta \mathbf{S} = (0, \delta A_y \cos Qy, -\delta A_z \sin Qy)$$

Susceptibilities

$$\epsilon_{ij} = \begin{pmatrix} \epsilon_{xx} & 0 & 0 \\ 0 & \epsilon_{yy} & 0 \\ 0 & 0 & \epsilon_{zz} \end{pmatrix}$$

$$\mu_{ij} = \begin{pmatrix} \mu_{xx} & 0 & 0 \\ 0 & \mu_{yy} & 0 \\ 0 & 0 & \mu_{zz} \end{pmatrix}$$

$$\alpha_{ik}^{em} = \begin{pmatrix} 0 & \alpha_{xy}^{em} & 0 \\ \alpha_{yx}^{em} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\alpha_{ik}^{me} = \begin{pmatrix} 0 & \alpha_{xy}^{me} & 0 \\ \alpha_{yx}^{me} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Time reversal invariance

$$\alpha_{ik}^{me} = -\alpha_{ki}^{em}$$

Maxwell equations

$$\mathbf{b} = +\mathbf{n} \times \mathbf{e}$$

$$\mathbf{d} = -\mathbf{n} \times \mathbf{h}$$

$$\mathbf{k} = \frac{\omega}{c} \mathbf{n}$$

$$d_i = \varepsilon_{ik} e_k + \alpha_{ik}^{em} h_k$$

$$b_i = \alpha_{ik}^{me} e_k + \mu_{ik} h_k$$

Fresnel equation

$$\left(\frac{n^2}{\varepsilon(\varphi)\mu_{zz}} - 1 \right) \left(\frac{n^2}{\mu(\varphi)\varepsilon_{zz}} - 1 \right)$$

$$+ \frac{\alpha^2}{\varepsilon_{xx}\mu_{yy}} \left(\frac{n^2 \cos^2 \varphi}{\varepsilon_{yy}\mu_{zz}} - 1 \right) \left(\frac{n^2 \sin^2 \varphi}{\varepsilon_{zz}\mu_{xx}} - 1 \right) = 0$$

$$\frac{1}{\varepsilon(\varphi)} = \frac{\sin^2 \varphi}{\varepsilon_{xx}} + \frac{\cos^2 \varphi}{\varepsilon_{yy}}$$

$$\frac{1}{\mu(\varphi)} = \frac{\sin^2 \varphi}{\mu_{xx}} + \frac{\cos^2 \varphi}{\mu_{yy}}$$

$$\alpha = \alpha_{xy}^{em} = -\alpha_{yx}^{me}$$

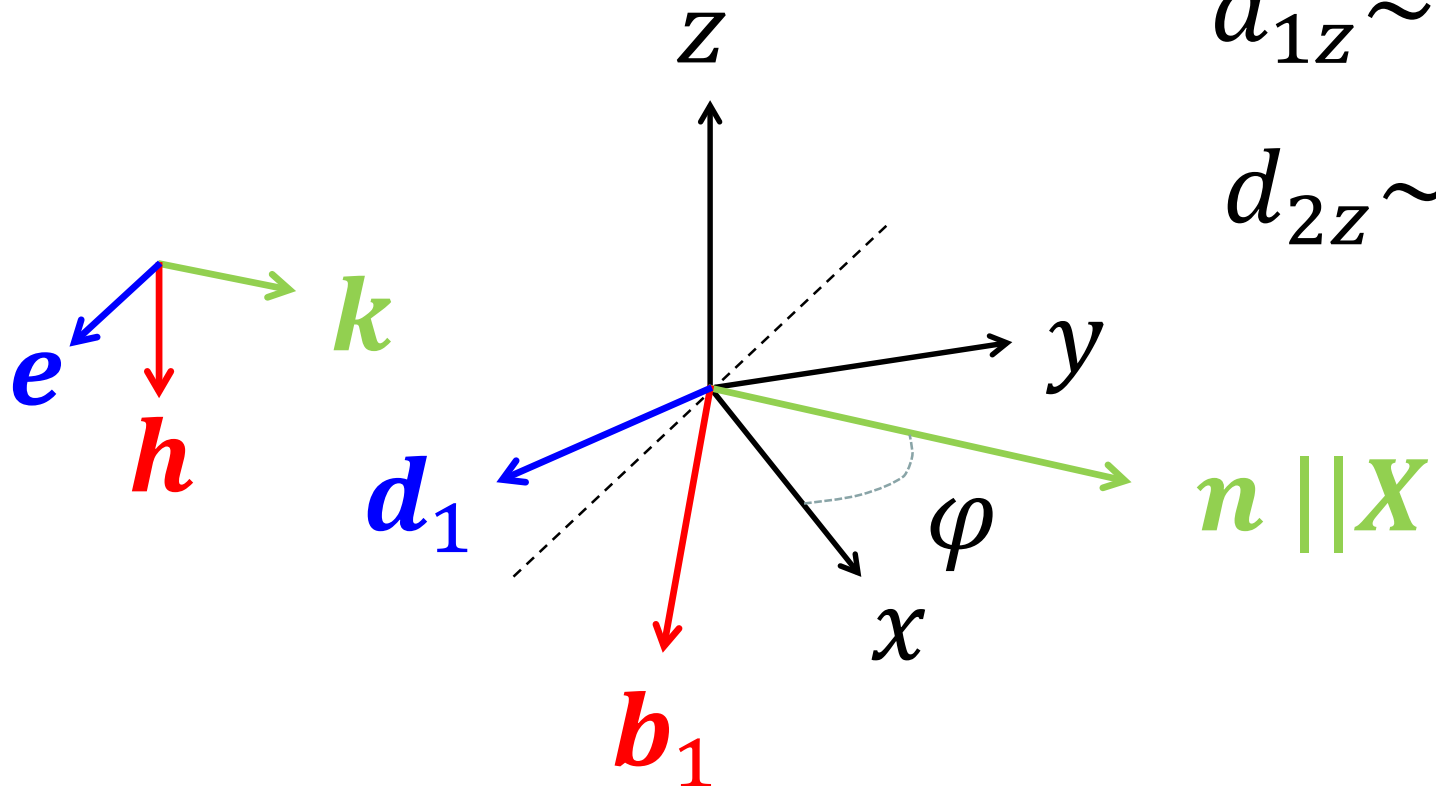
$$\alpha_{yx}^{em} = -\alpha_{xy}^{me} = 0$$

d_z -oscillations due to magnetoelectric coupling

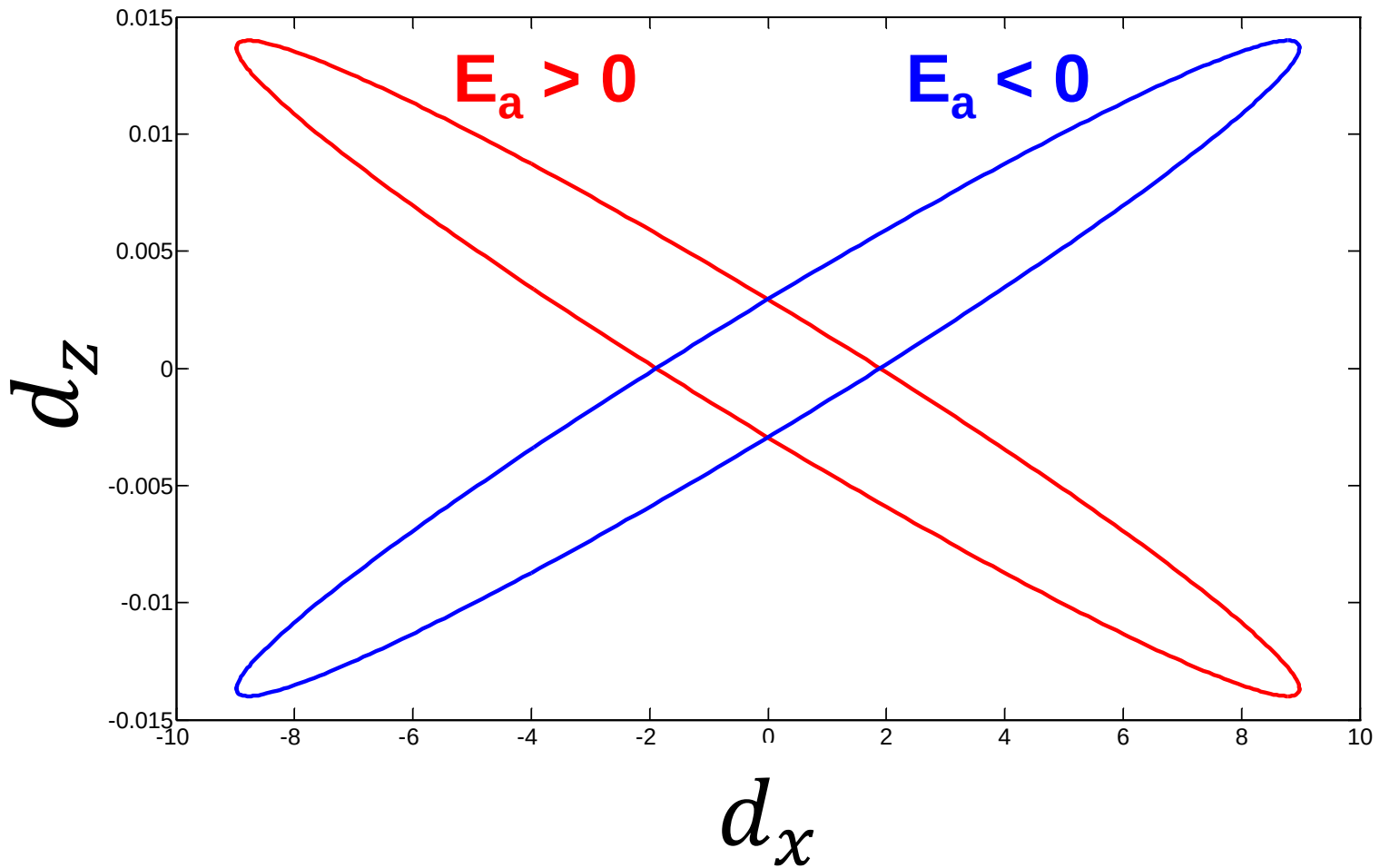
$$d_z = e^{-i\omega t} \left(d_{1z} e^{i\frac{\omega}{c}n_1 X} + d_{2z} e^{i\frac{\omega}{c}n_2 X} \right)$$

$$d_{1z} \sim \alpha$$

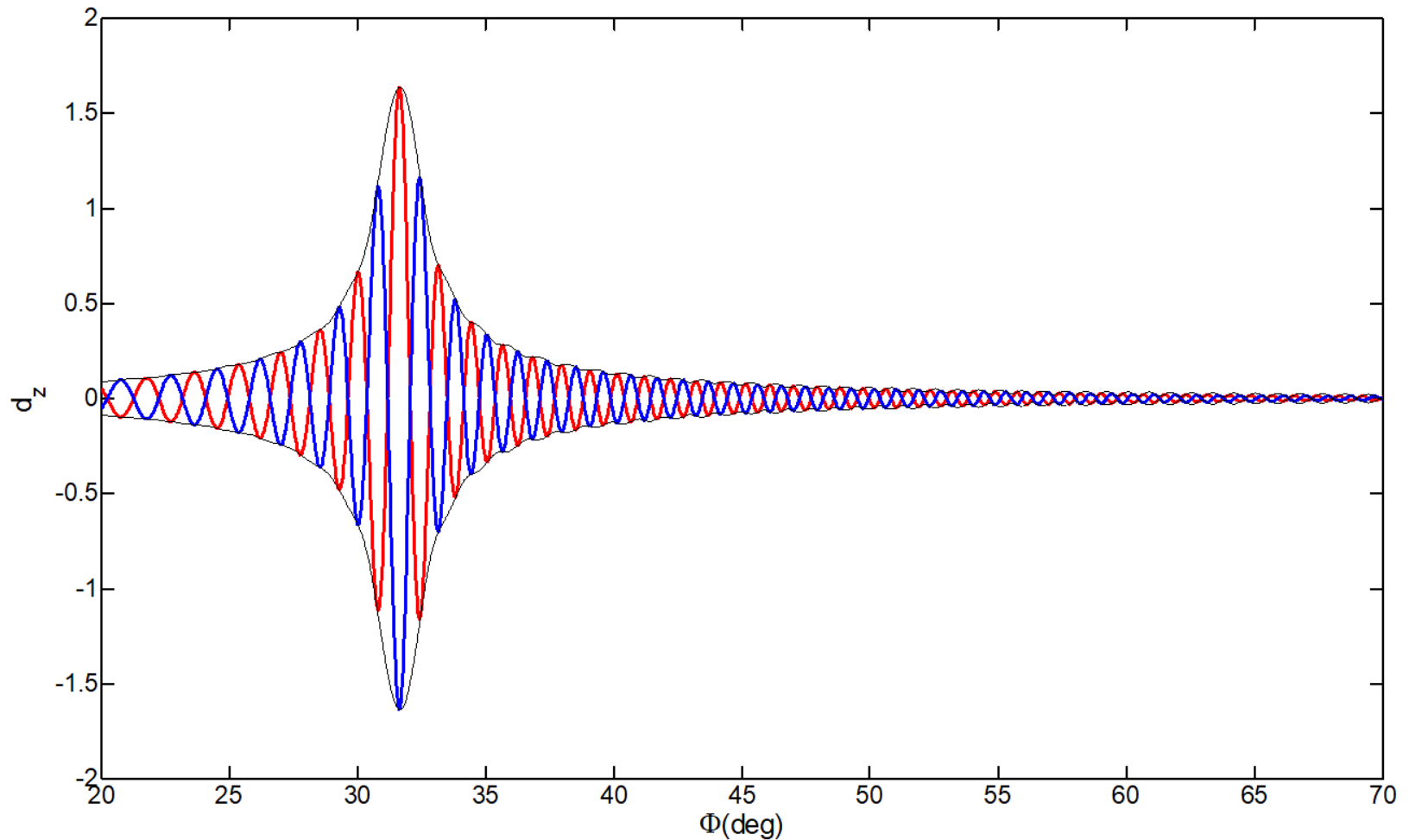
$$d_{2z} \sim \alpha$$



Elliptical polarization

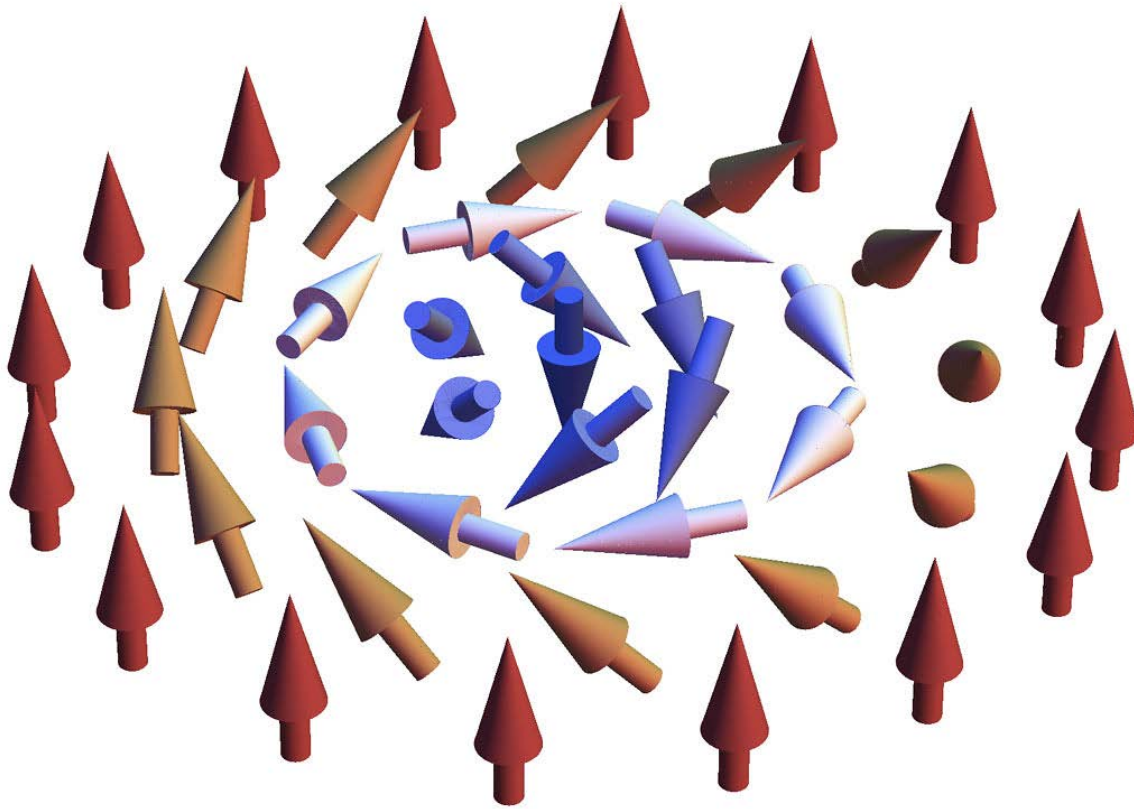


Enhancement near optical bi-axis



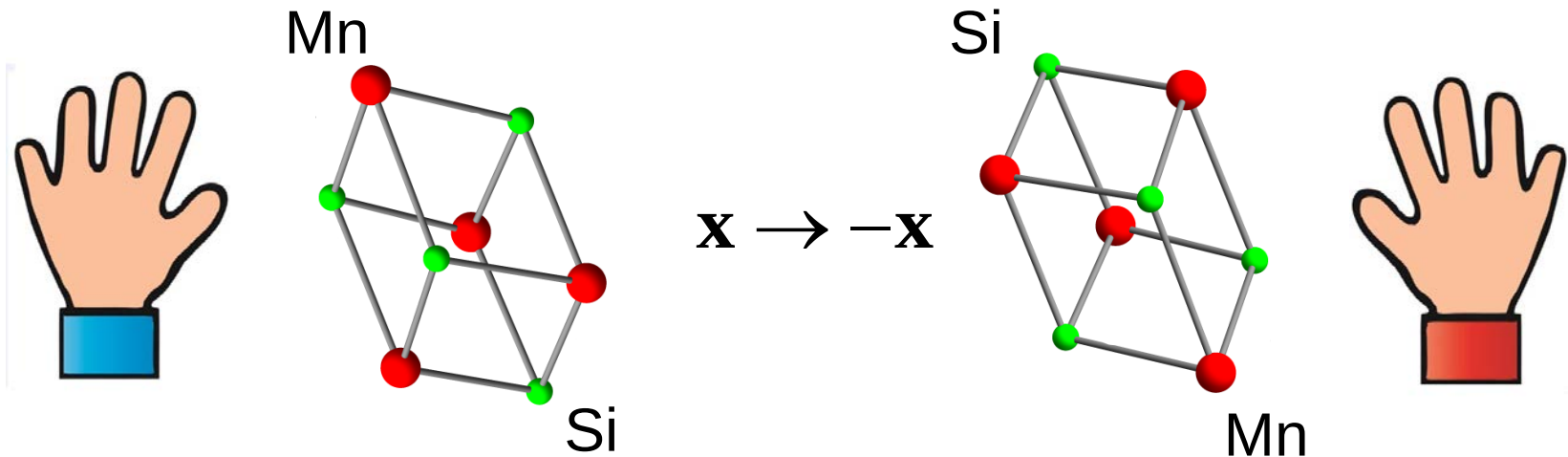
Magnetic skyrmion

$$n_x^2 + n_y^2 + n_z^2 = 1$$

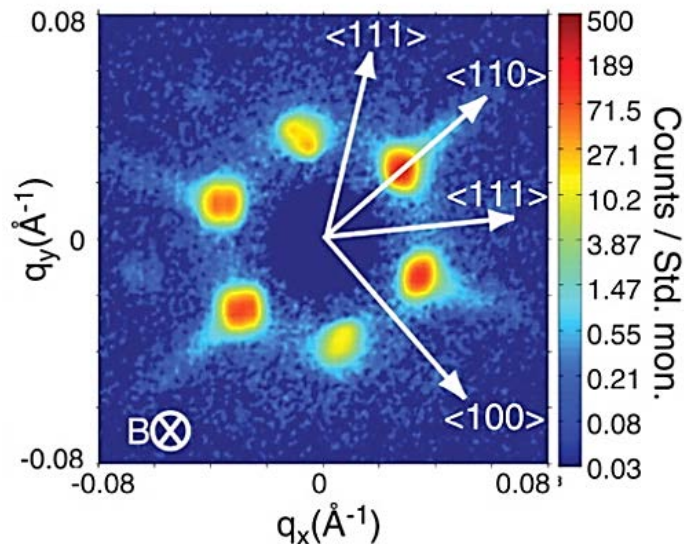


$$S_2 \rightarrow S_2$$

Skymions in chiral magnets

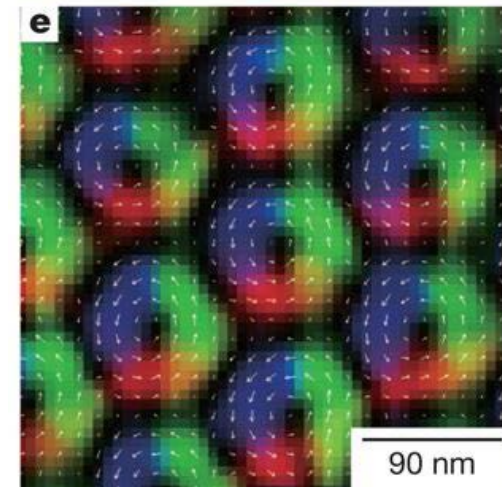


Neutron scattering



Mühlbauer et al, Science (2009)

Lorentz microscopy



X. Yu et al. Nature (2010)

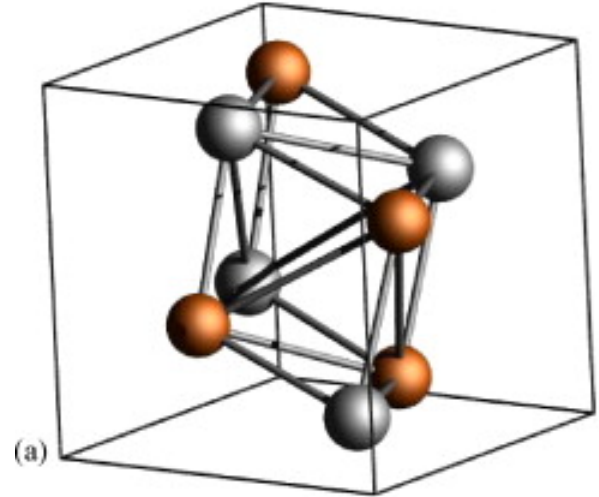
Non-centrosymmetric magnets

I. Dzyaloshinskii, Sov. Phys. JETP 19, 960 (1964)

P. Bak & M.H. Jensen, J. Phys. C: Solid State 13, L881 (1980)

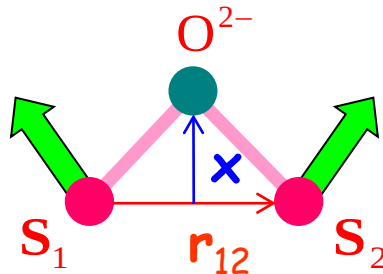
Cubic B20 structure $P32_1$

MnSi, $\text{Fe}_{1-x}\text{Co}_x\text{Si}$, FeGe



$$f = \frac{J}{2} (\nabla \mathbf{n})^2 + D \mathbf{n} \cdot [\nabla \times \mathbf{n}]$$

Exchange
energy

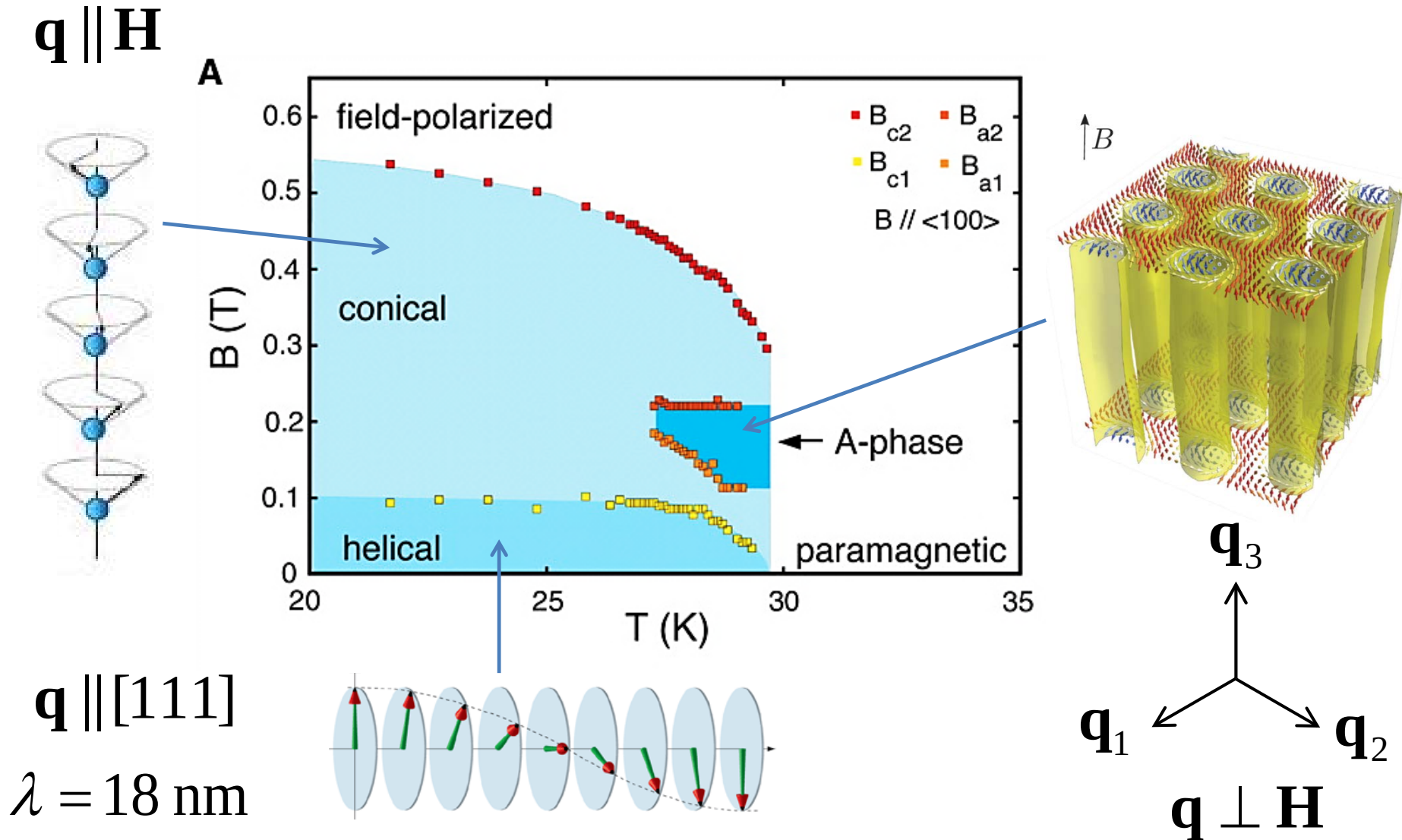


DM energy

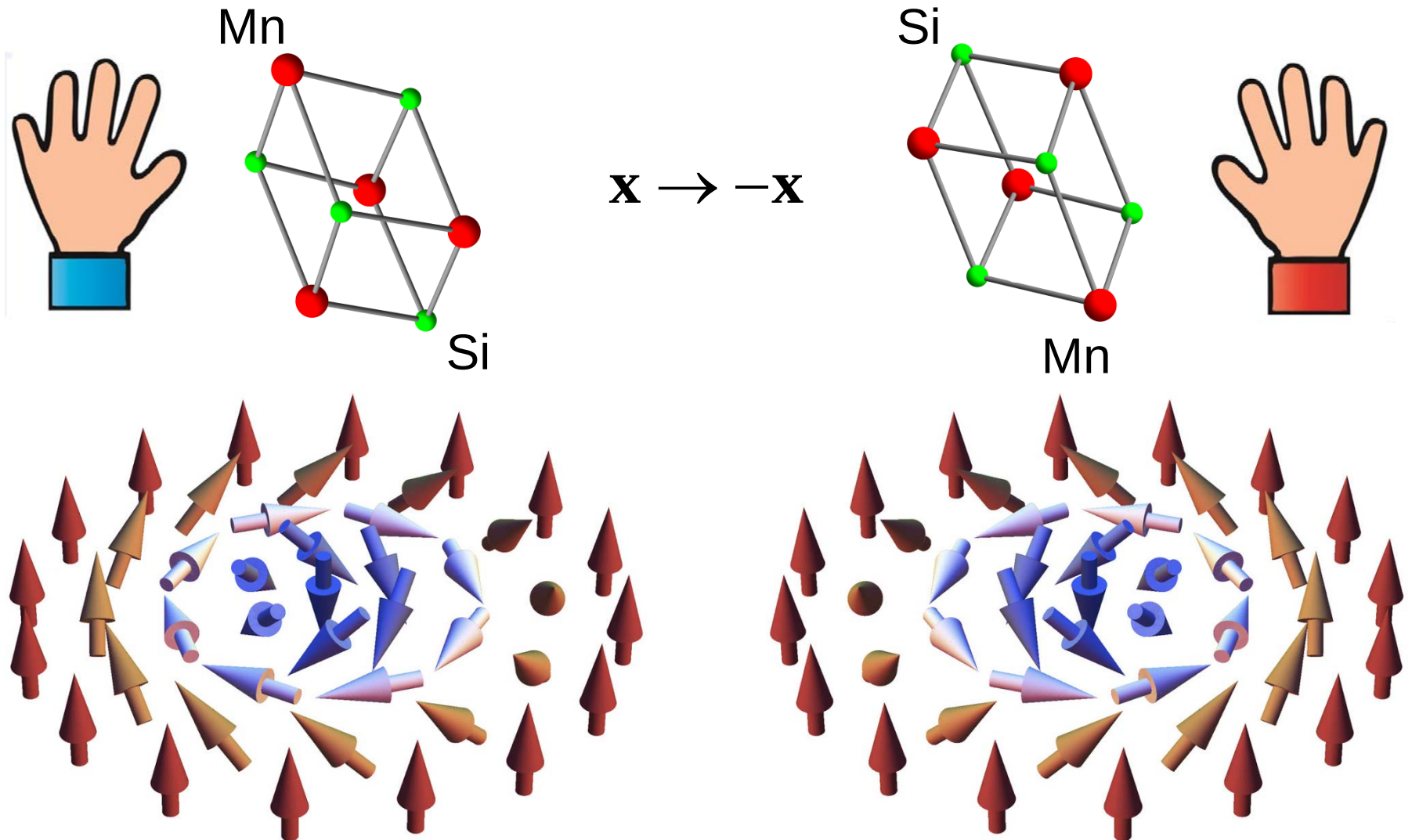
$$E_{DM} = \mathbf{D}_{12} \cdot [\mathbf{S}_1 \times \mathbf{S}_2]$$

Phase Diagram of MnSi

Mühlbauer et al, *Science* **323**, 915 (2009)

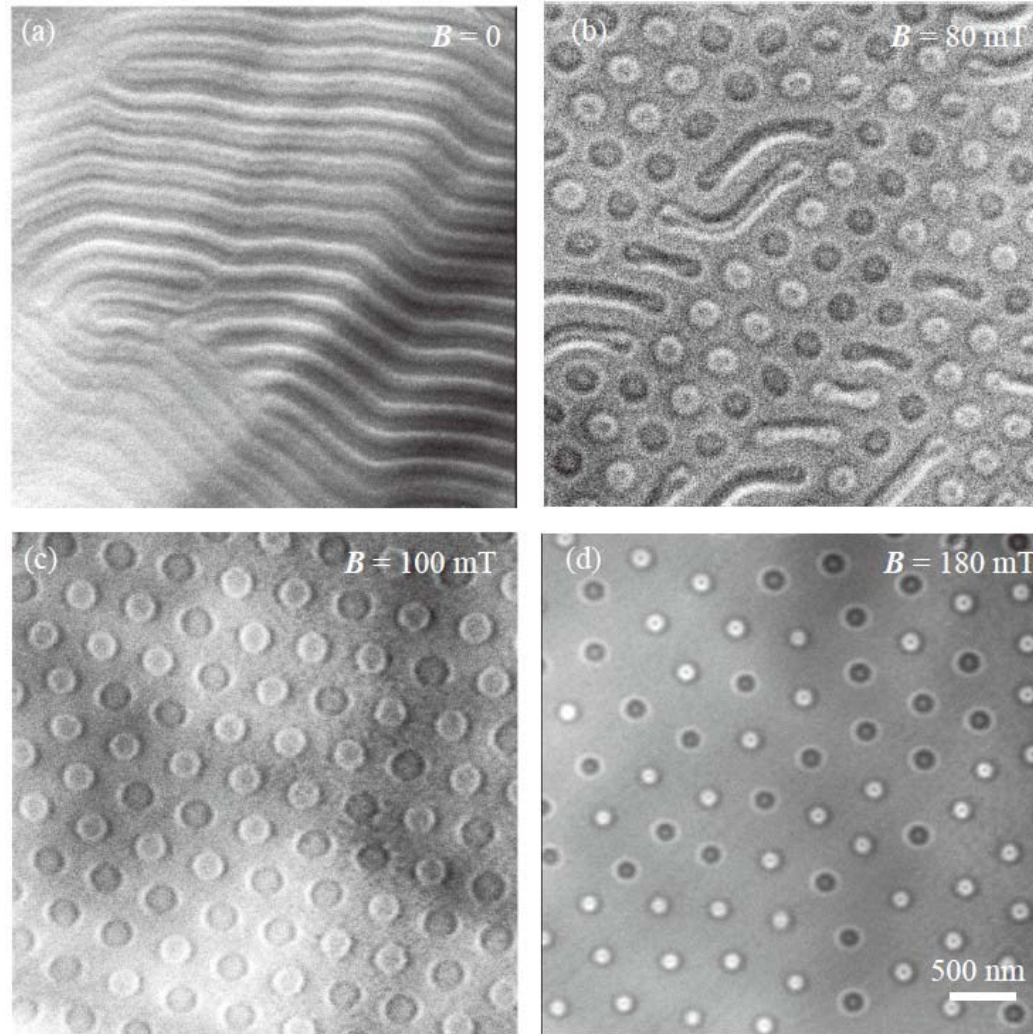


Skyrmion inversion



Topological charge does not change under inversion

Stripes & magnetic bubbles



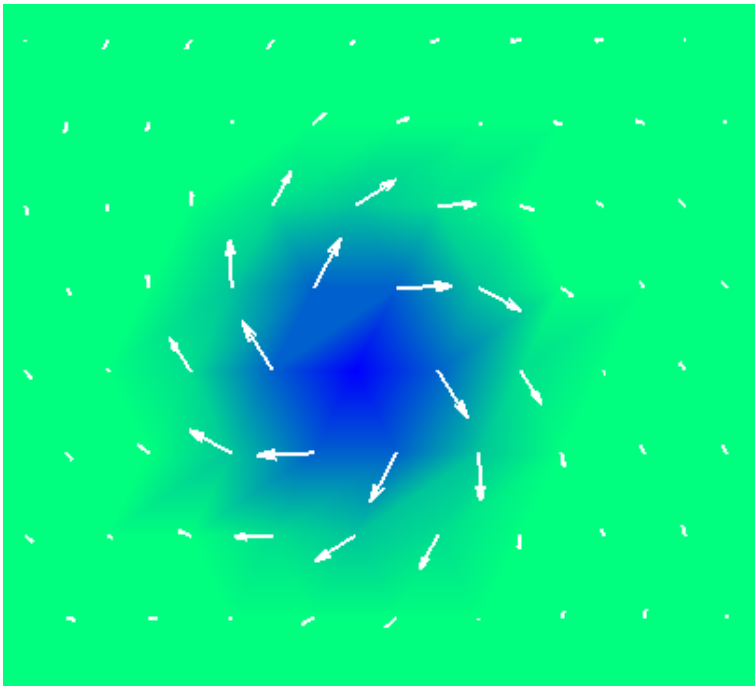
X. Yu et al. PNAS 109, 8856 (2012)

Skyrmions in frustrated magnets

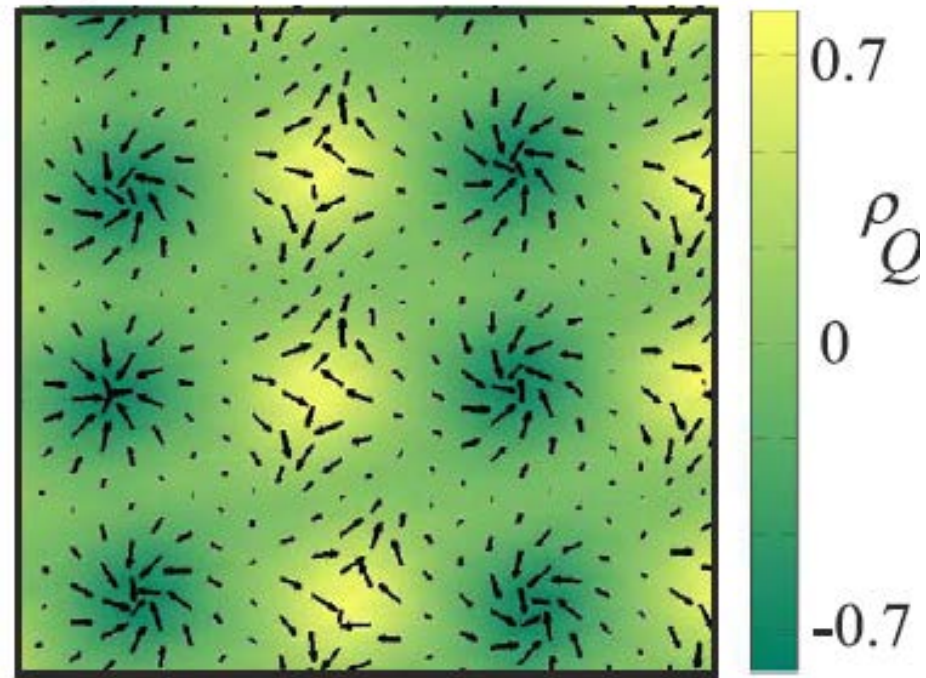
Arbitrary helicity and vorticity

$$\Theta = \Theta(\rho) \quad \Phi = v\varphi + \chi$$

$$Q = v \frac{(n_z(0) - n_z(\infty))}{2}$$

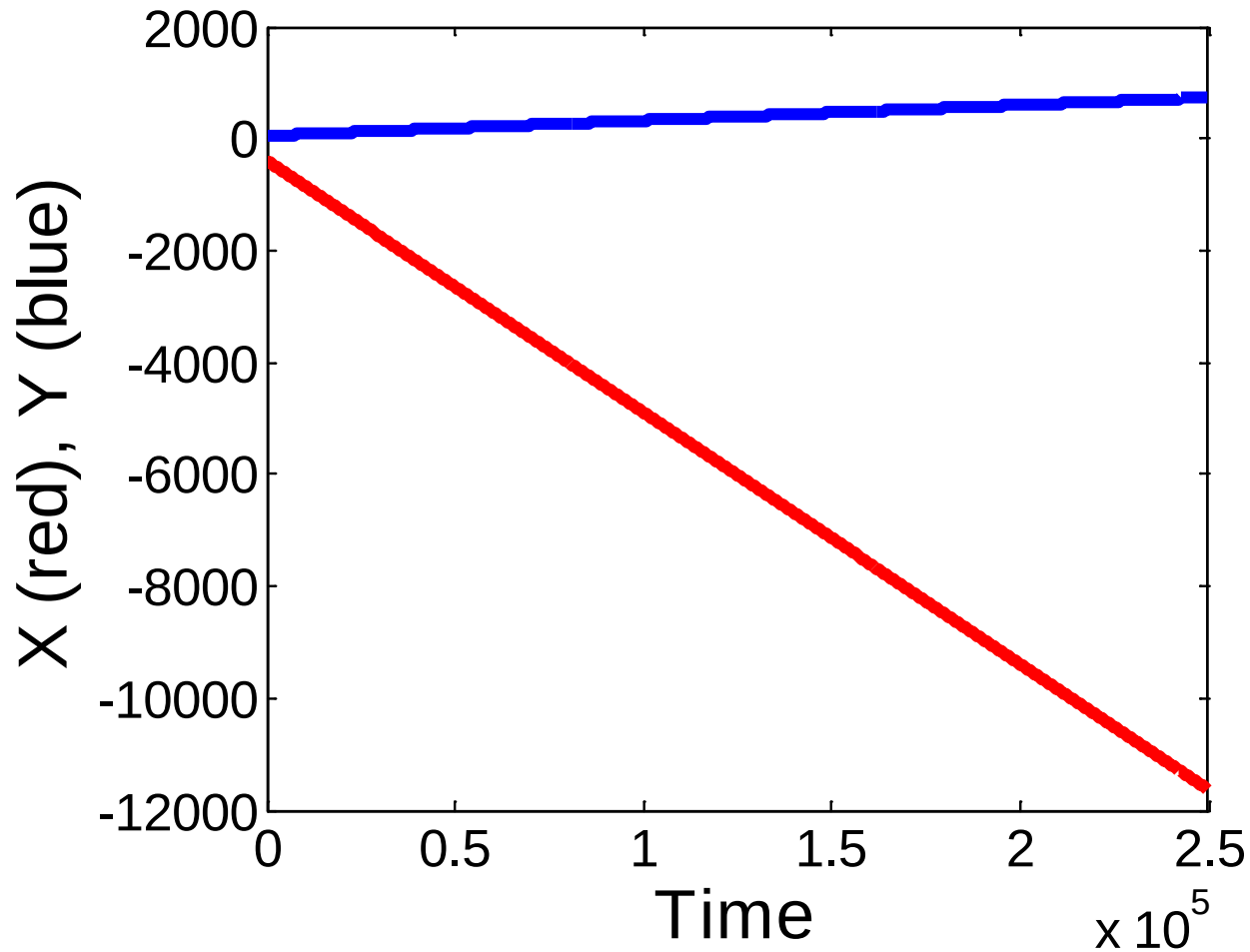


$$\mathbf{S} \rightarrow R_z(\chi) \mathbf{S}$$



$$S_x \rightarrow -S_x$$

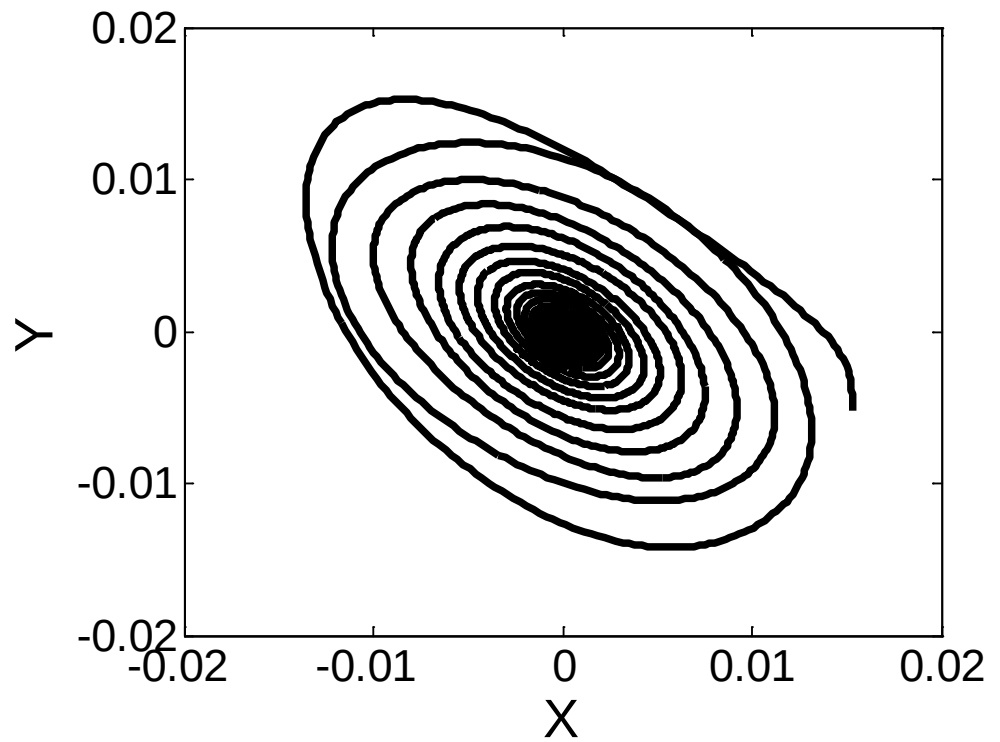
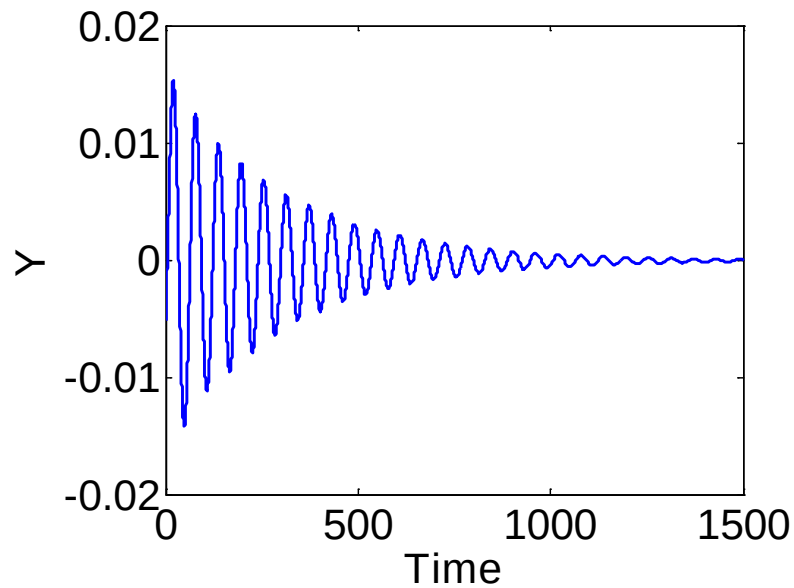
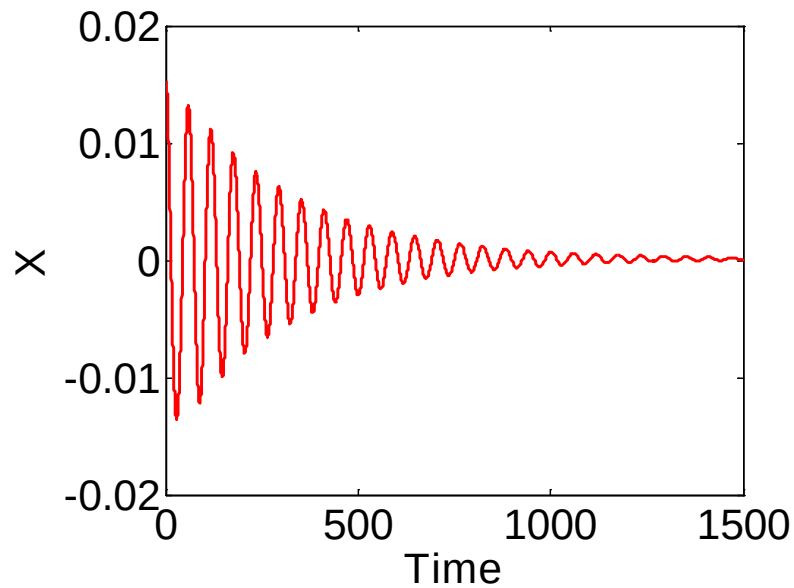
Current-driven skyrmion: translational motion



$$j_x = 0.05$$

Does skyrmion have
a mass?

Coordinate relaxation



Relaxation of the massive skyrmion

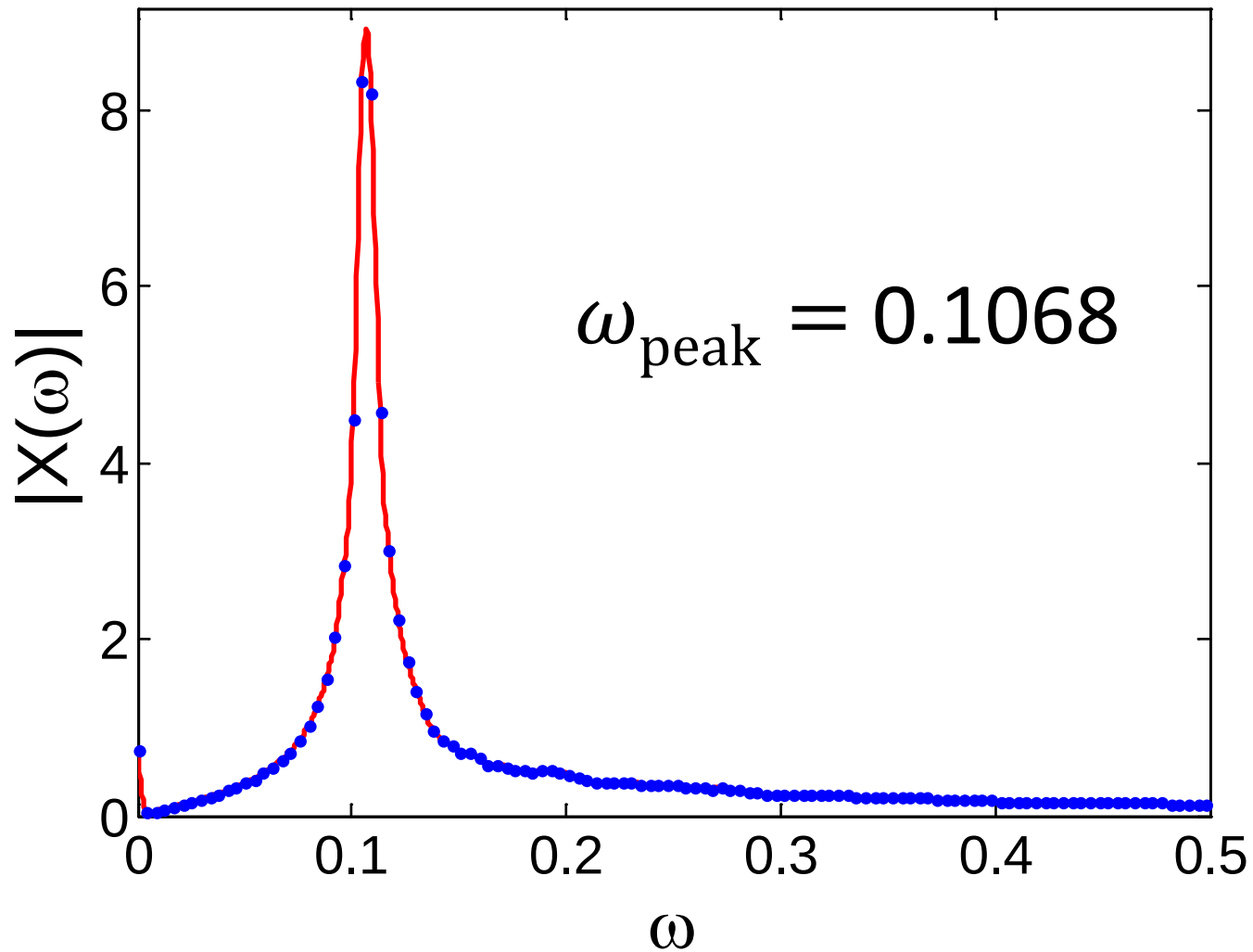
Center-of-mass coordinates:

$$\mathbf{m}(\mathbf{x}, t) = \mathbf{m}(\mathbf{x} - \mathbf{R}(t)) \quad \mathbf{R} = (X, Y)$$

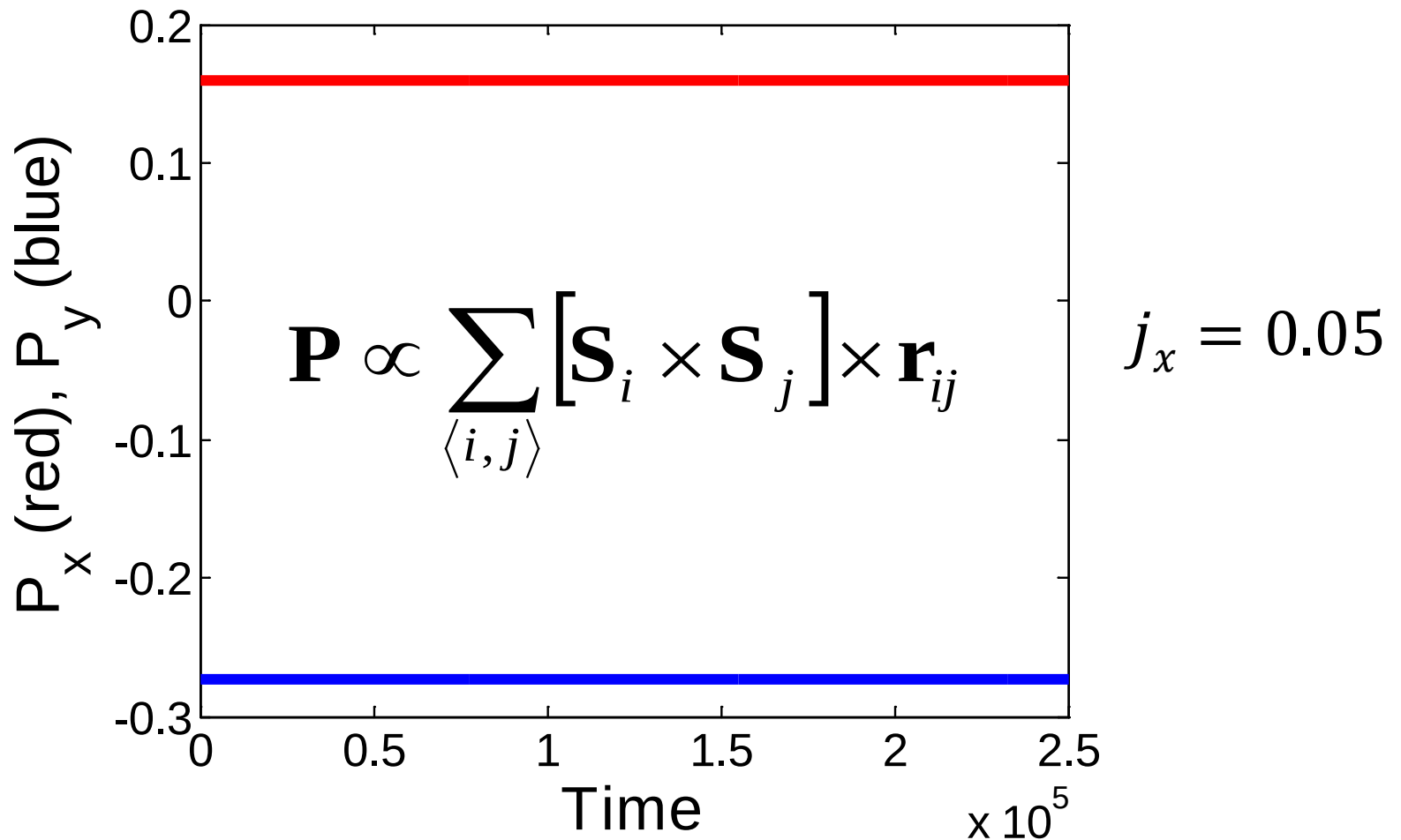
$$\begin{cases} M \ddot{X} + \alpha \Gamma \dot{X} - G \dot{Y} = 0 \\ M \ddot{Y} + G \dot{X} + \alpha \Gamma \dot{Y} = 0 \end{cases}$$

$$\omega_{\pm} = \frac{\pm |G| - i\alpha\Gamma}{M}$$

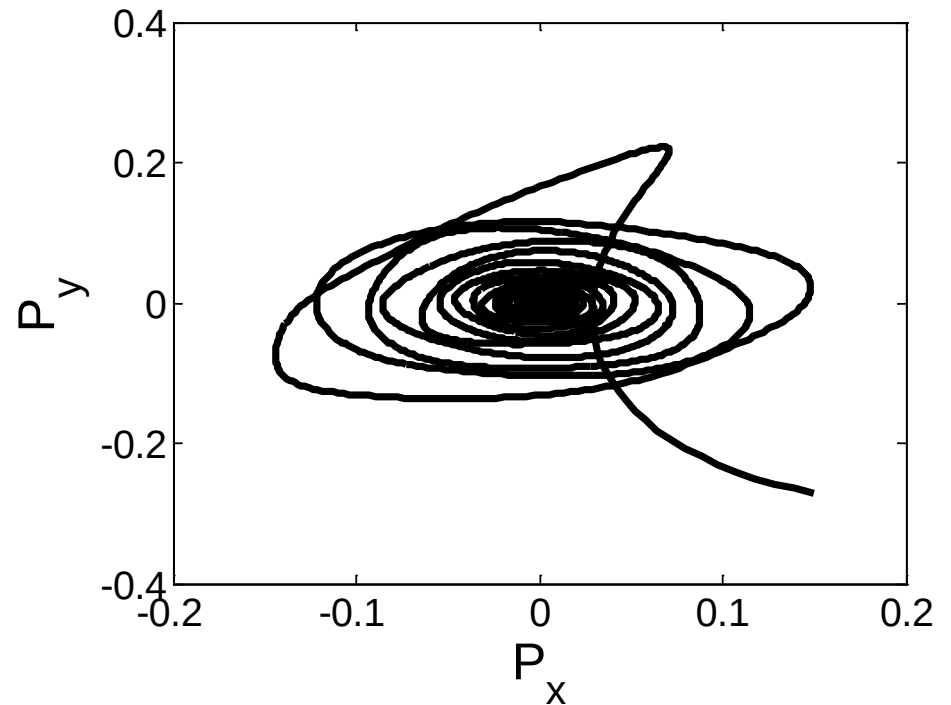
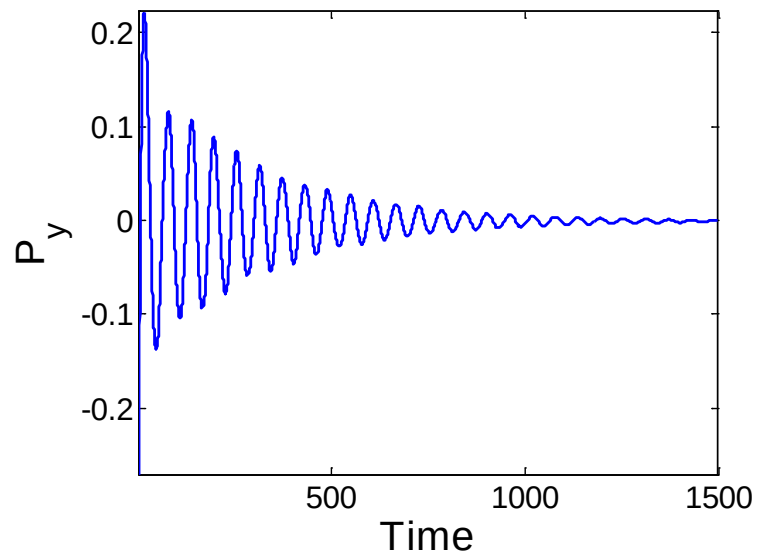
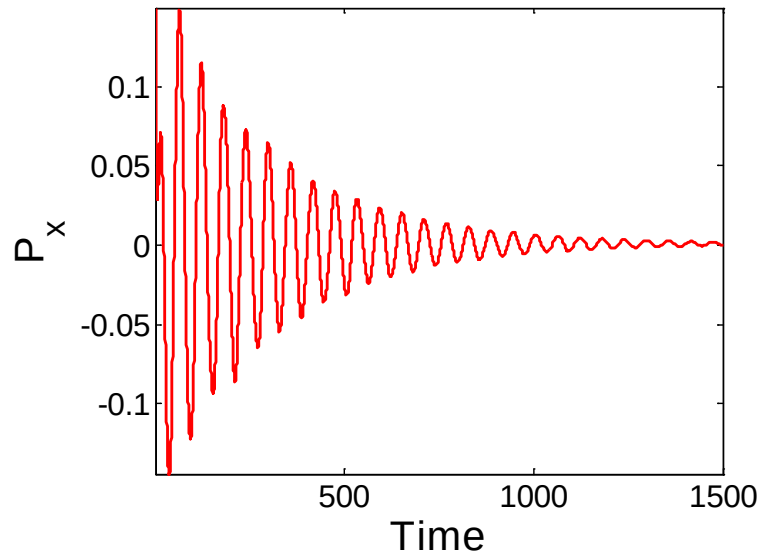
Fourier transform of $X(t)$



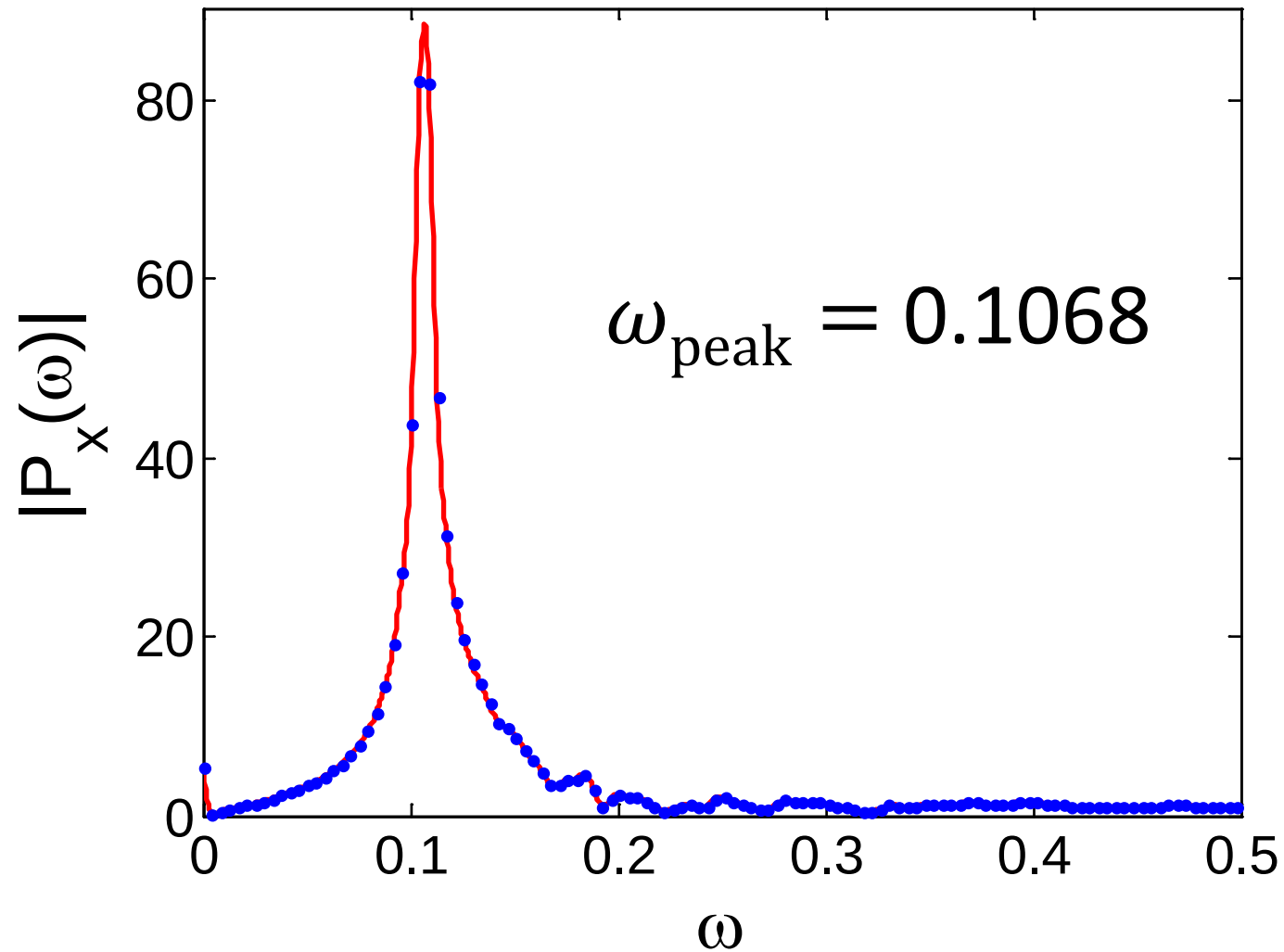
Current-driven skyrmion: electric polarization



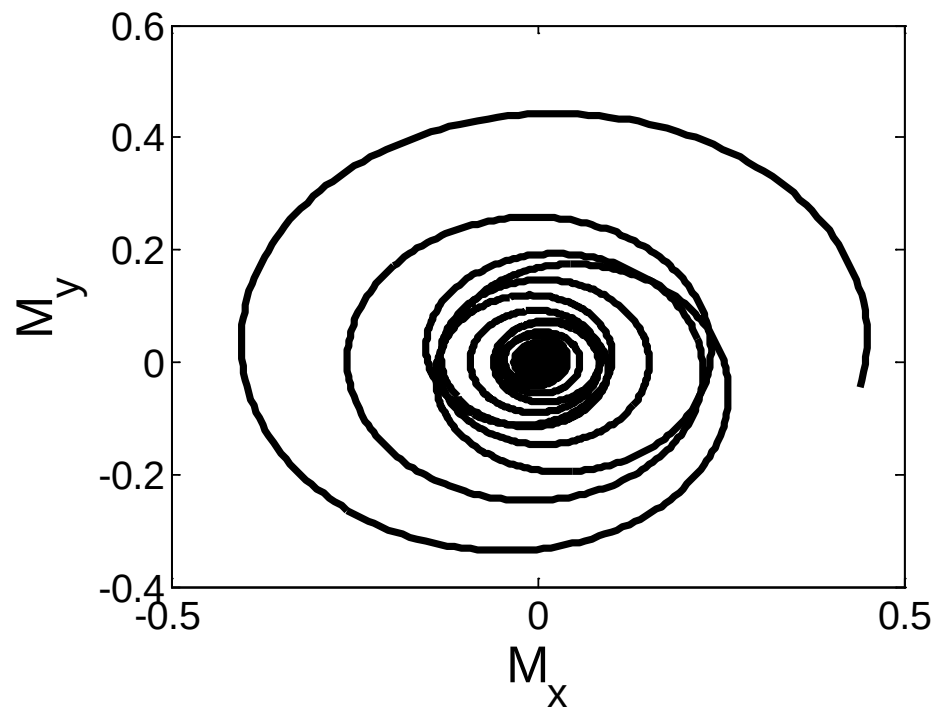
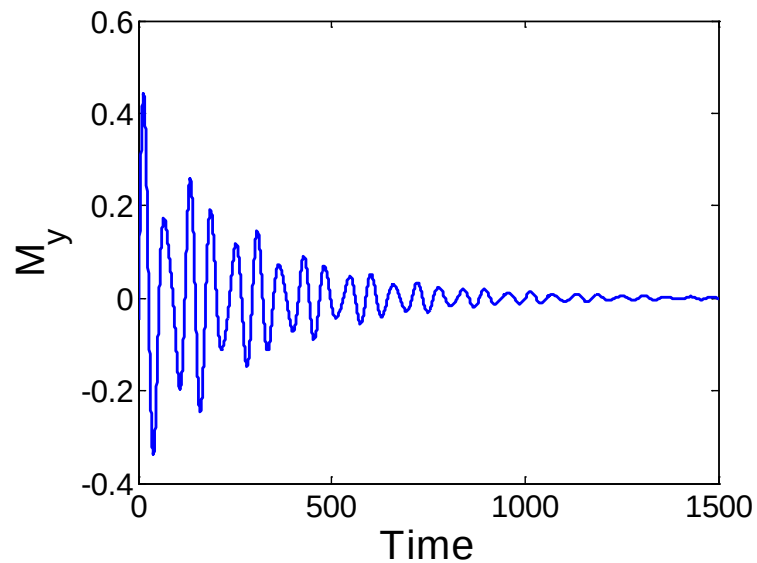
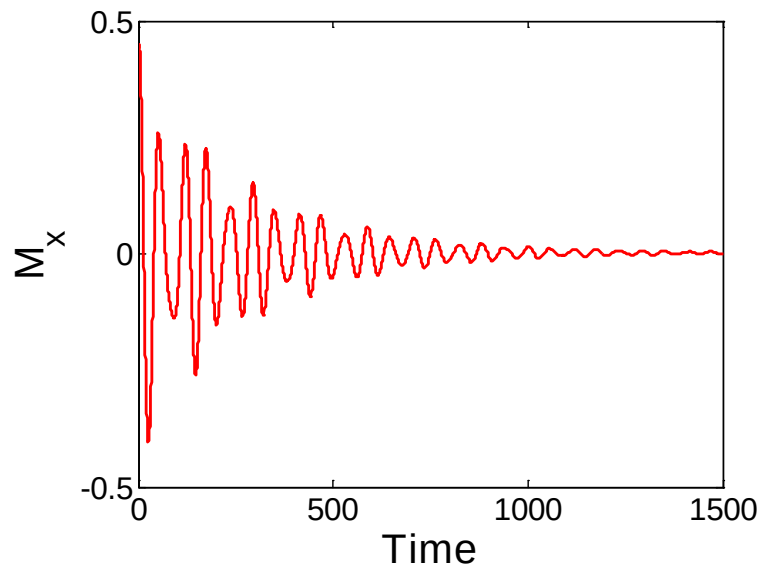
Electric polarization relaxation



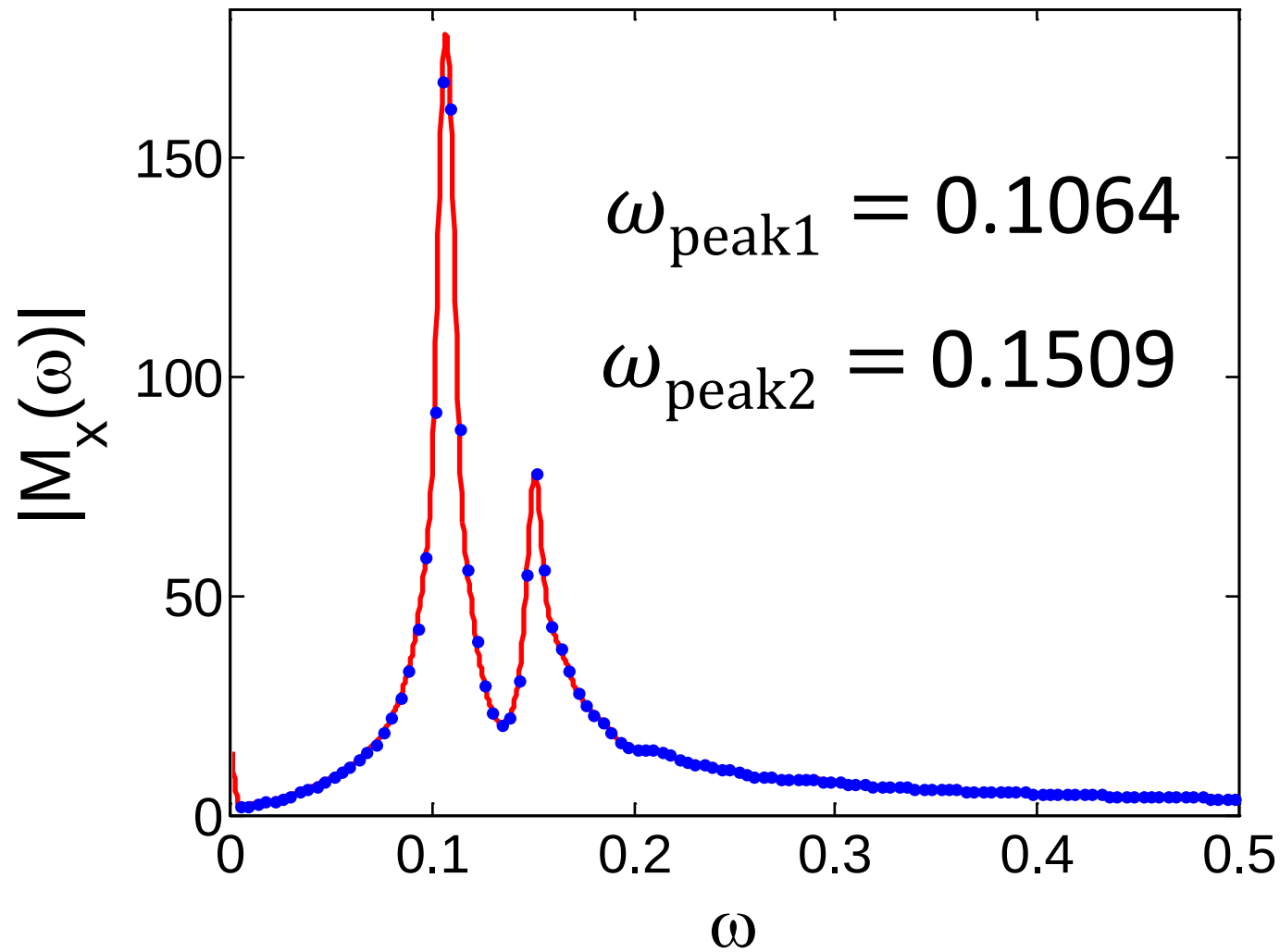
Fourier transform of $P_x(t)$



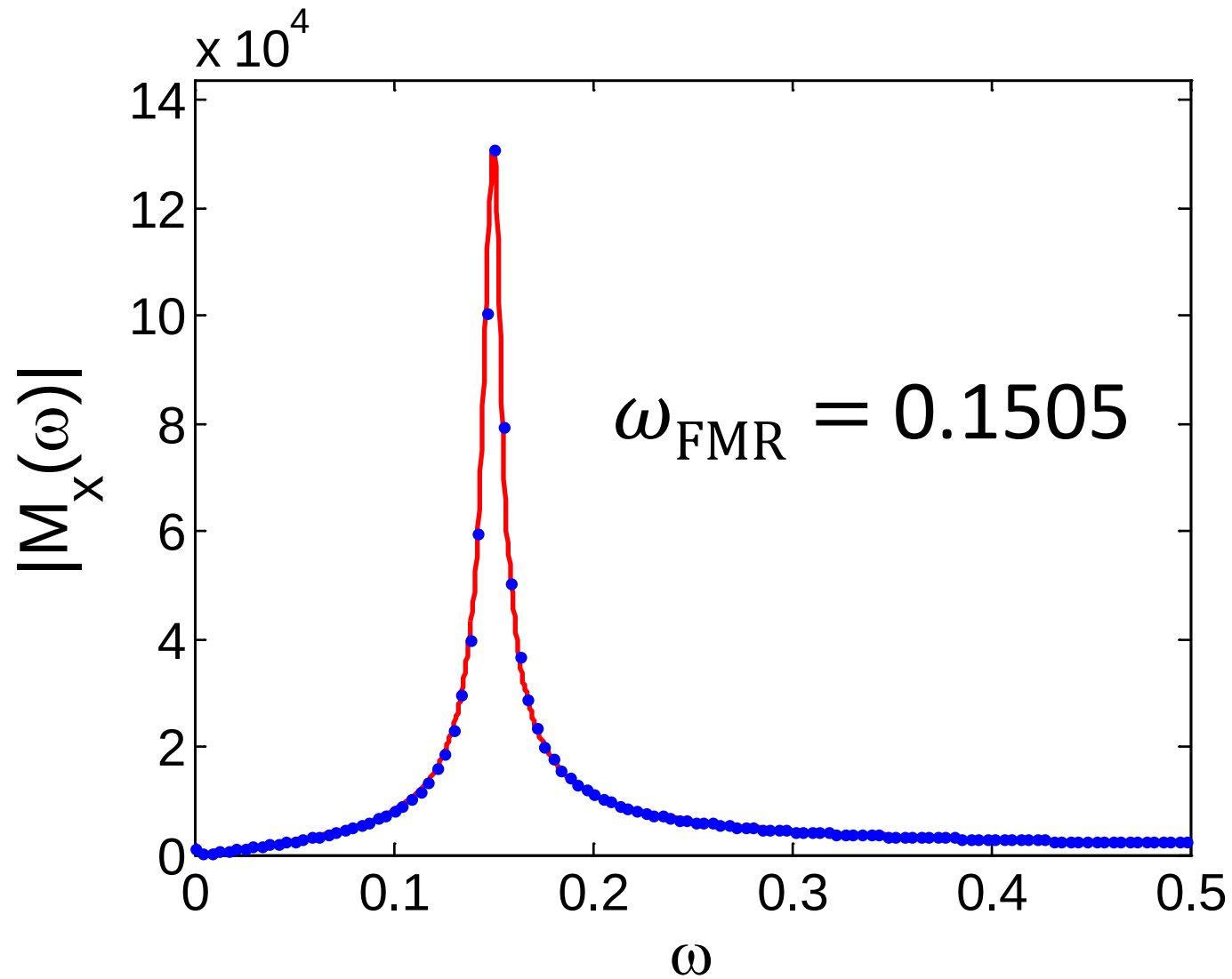
Magnetization relaxation



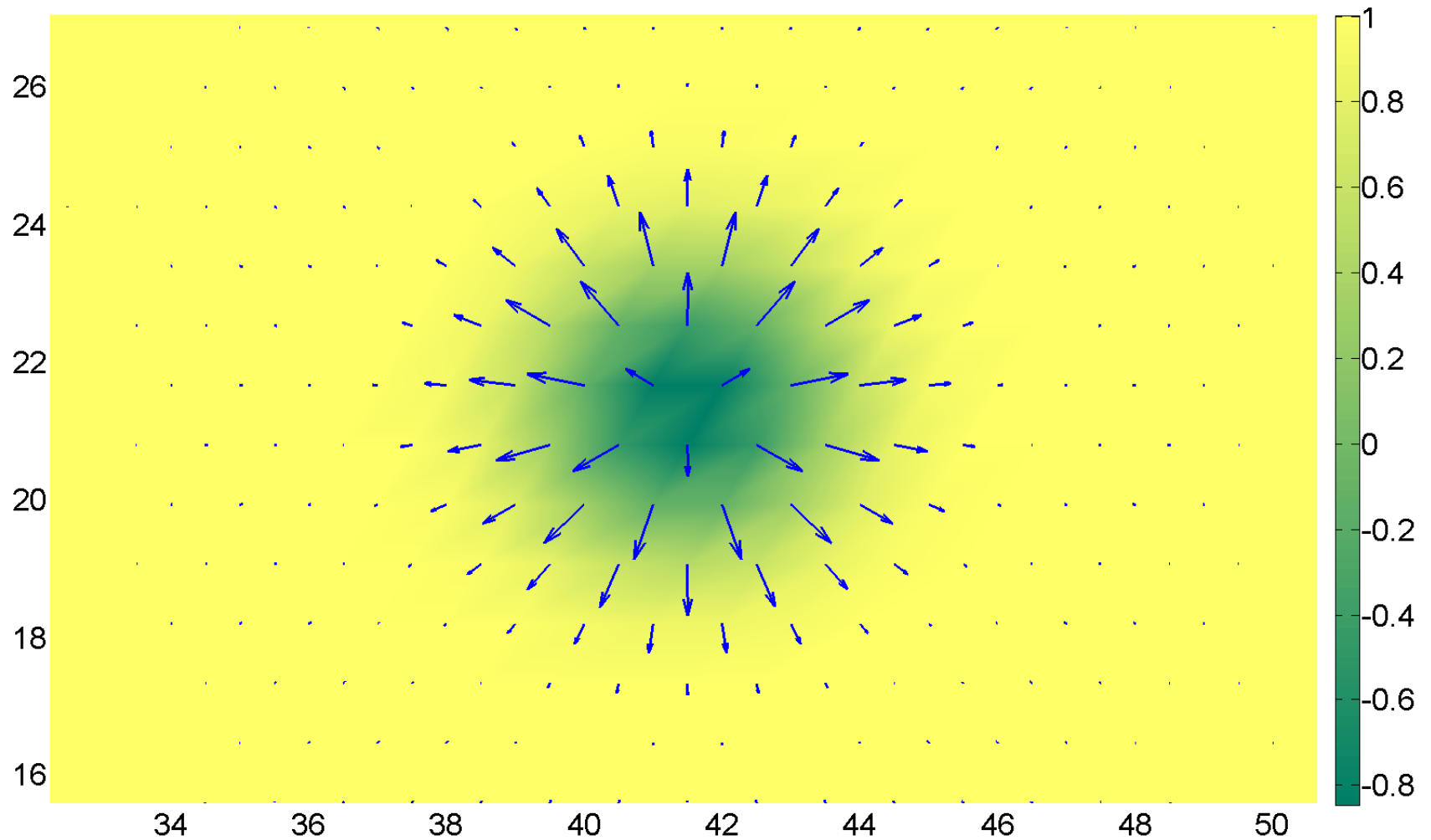
Fourier transform of $M_x(t)$



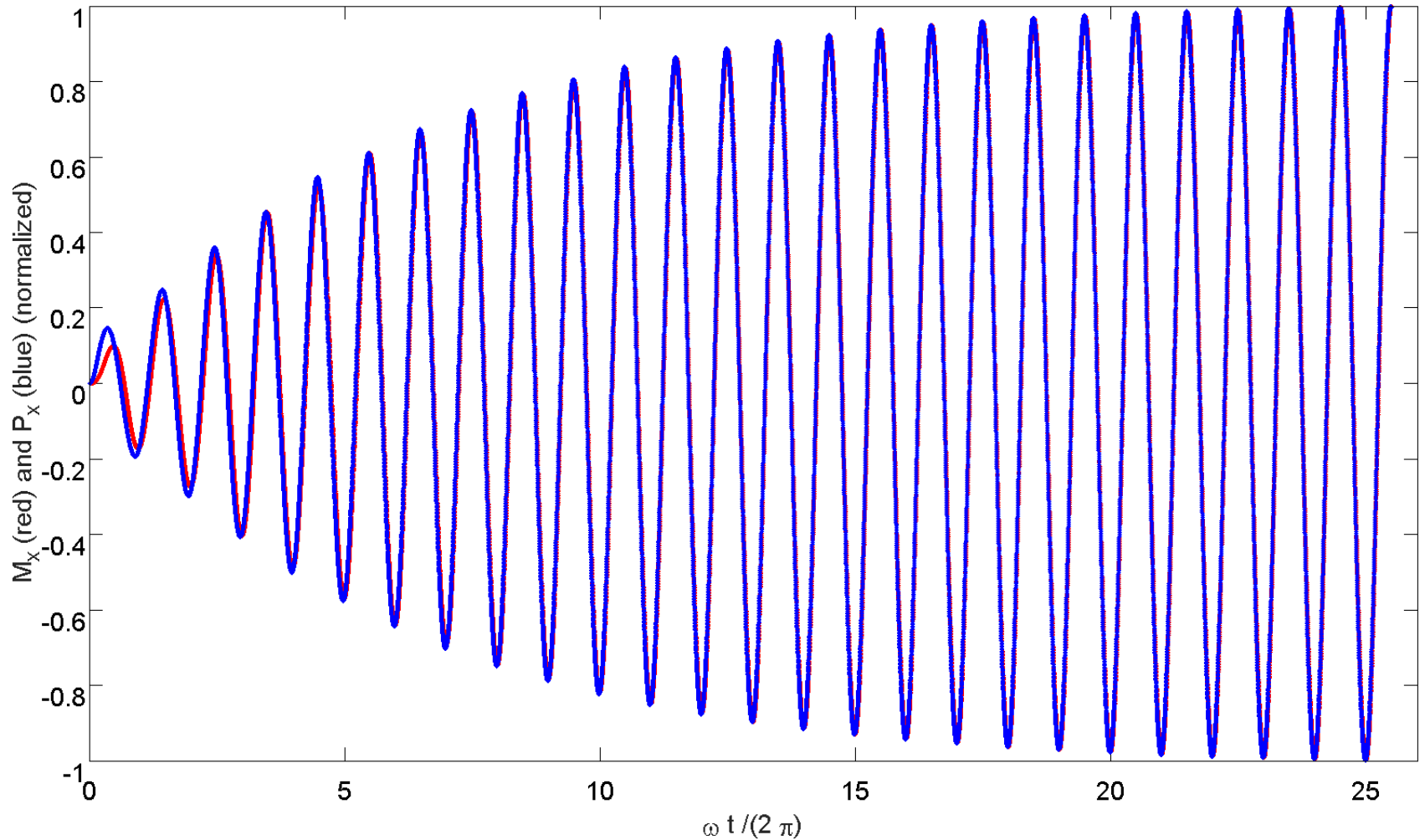
Bulk ferromagnetic resonance



Néel Skyrmion

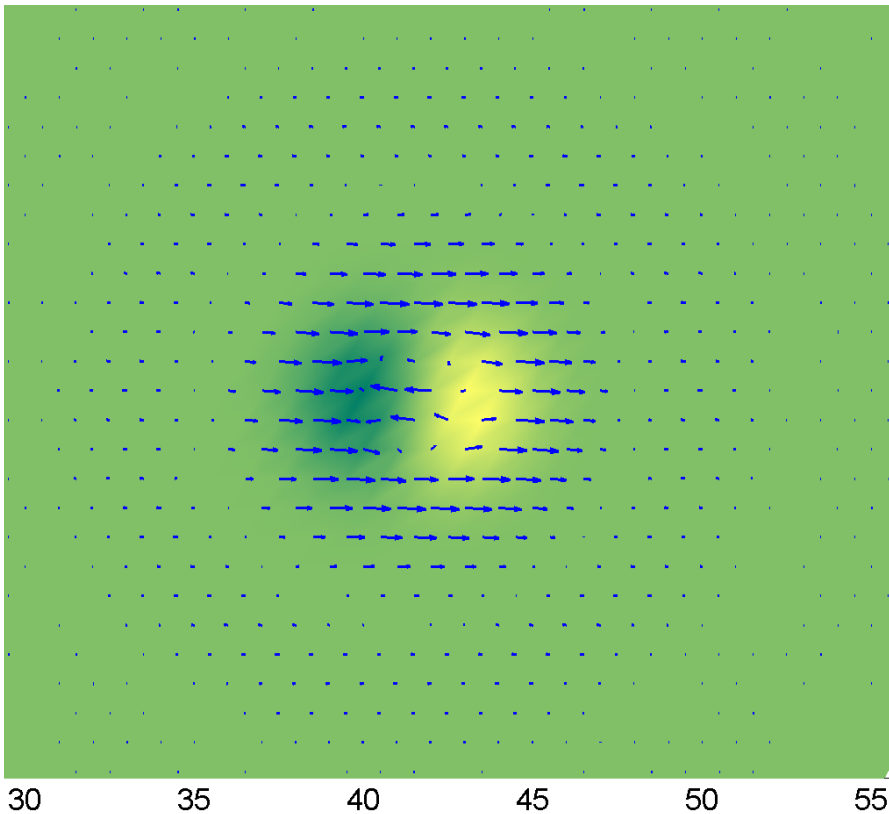


Excitation of magnetic moment with electric field



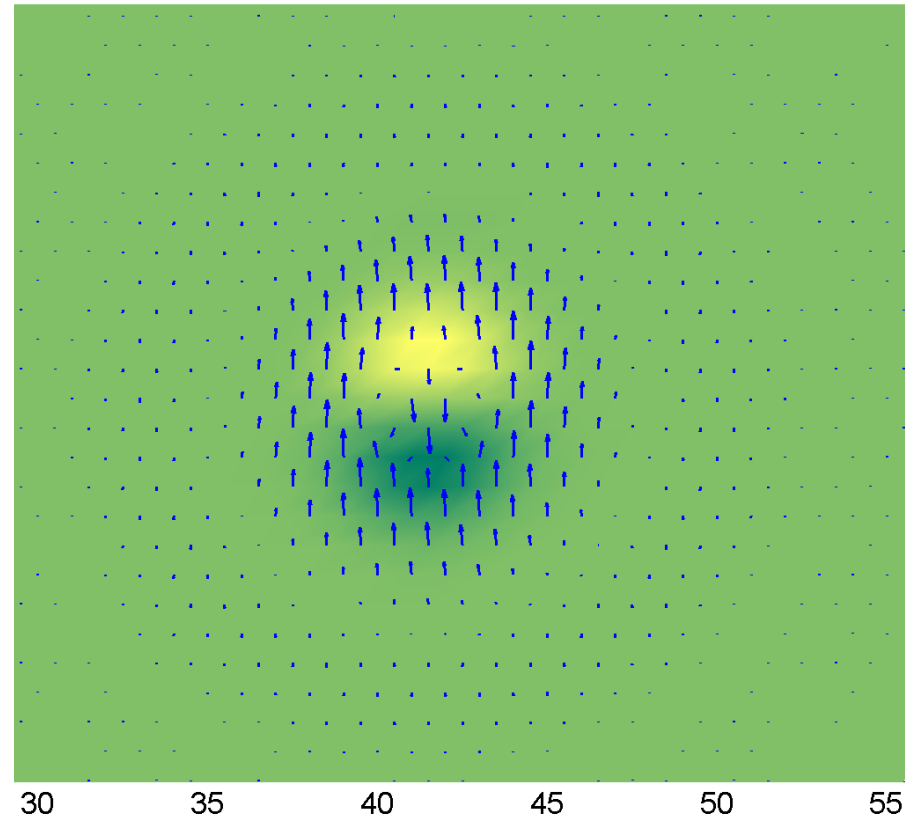
Localized mode with $\omega = 0.1068$

Q_1



P_x, M_x

Q_2



P_y, M_y

Coupling to one mode

$$\frac{(1 + \alpha^2)}{\omega_\nu} \ddot{X} + \alpha \Gamma_{\nu 2, \nu 2} \dot{X} + \dot{Y} = 0$$

$$\frac{(1 + \alpha^2)}{\omega_\nu} \ddot{Y} + \alpha \Gamma_{\nu 1, \nu 1} \dot{Y} - \dot{X} = 0$$

$$\Gamma_{\nu 1, \nu 1} = \Gamma_{\nu 2, \nu 2} \approx 1.23$$

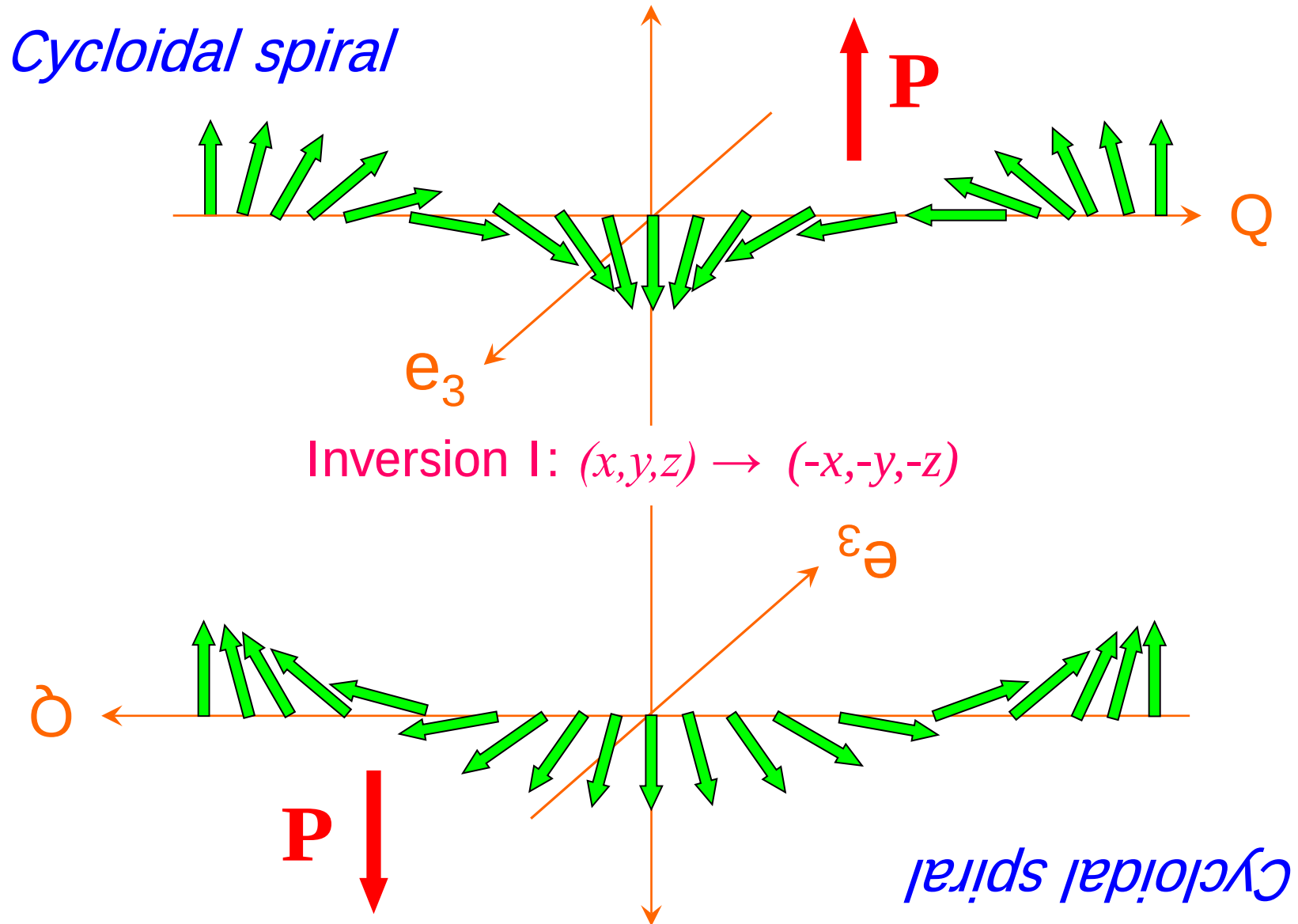
which has the same form as

$$\begin{cases} M \ddot{X} + \alpha \Gamma \dot{X} - G \dot{Y} = 0 \\ M \ddot{Y} + G \dot{X} + \alpha \Gamma \dot{Y} = 0 \end{cases}$$

$$M = (1 + \alpha^2) \frac{|G|}{\omega_\nu}$$

$$\frac{\Gamma}{|G|} \approx 1.25$$

Breaking of inversion symmetry by spin ordering



Skymion electric dipole moment

Magneto-electric coupling

$$f_{me} = -g \mathbf{E} \cdot [\mathbf{m}(\nabla \cdot \mathbf{m}) - (\mathbf{m} \cdot \nabla)\mathbf{m}]$$

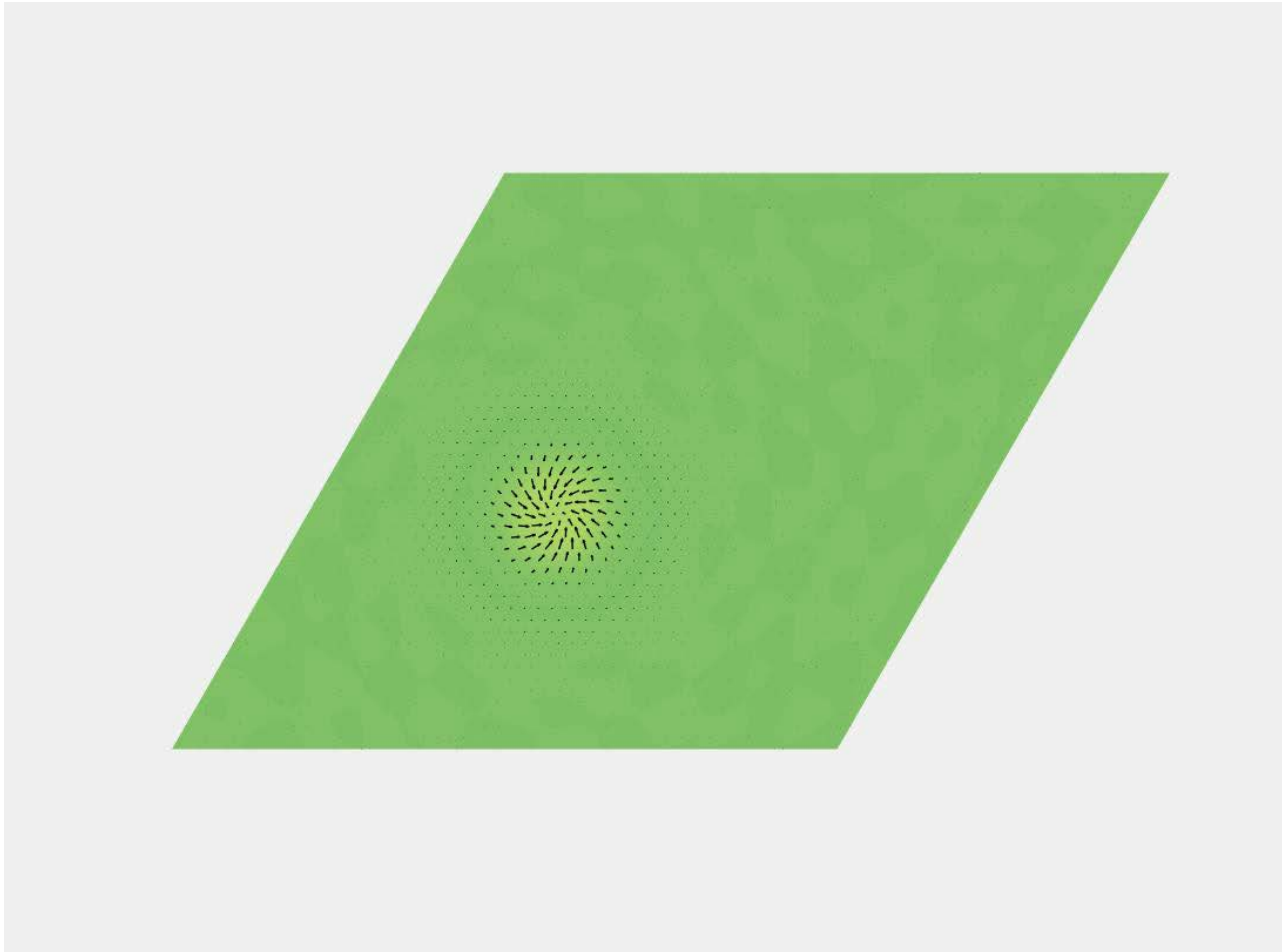
Polarization induced by skyrmion

$$P_z = -\frac{\partial f}{\partial E_z} = g \left[\frac{d\Theta}{dr} + v \frac{\sin 2\Theta}{2r} \right] \cos((v-1)\varphi + \chi)$$

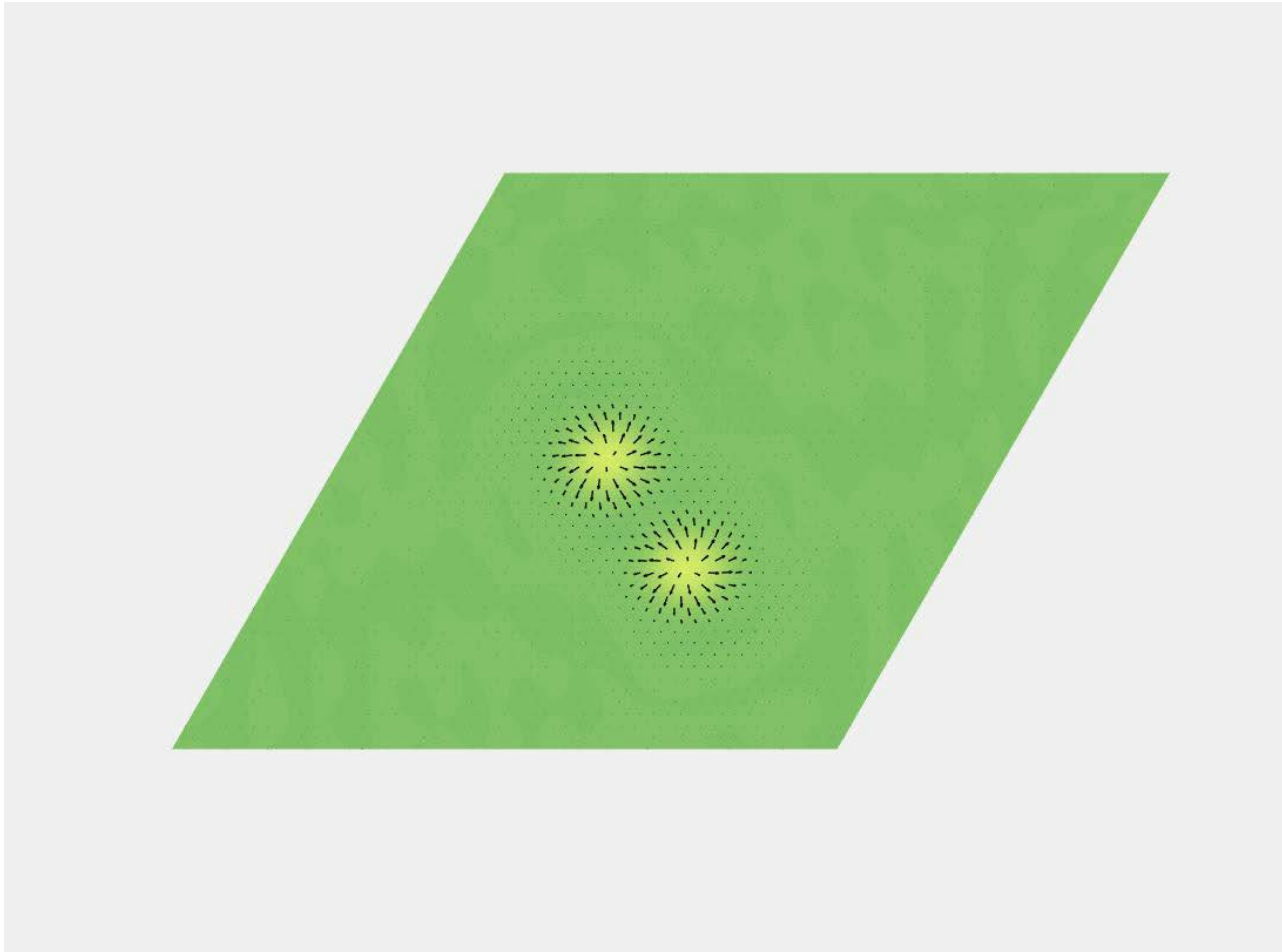
Dipole moment for vorticity $v = +1$

$$D_z(\chi) = D(0) \cos \chi$$

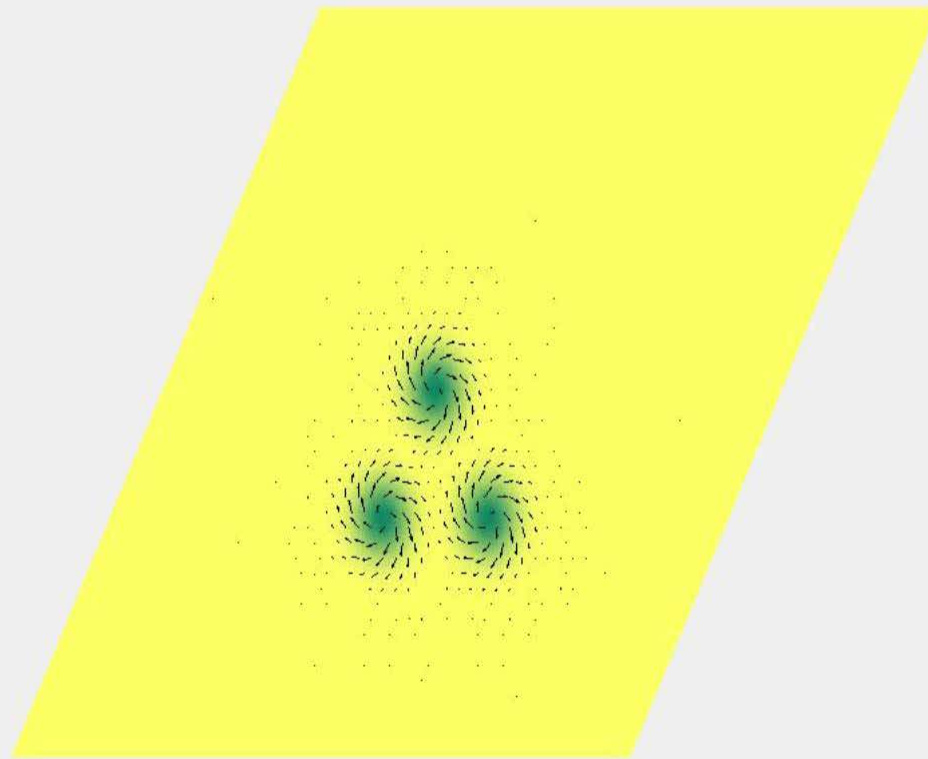
Electric excitation of skyrmion



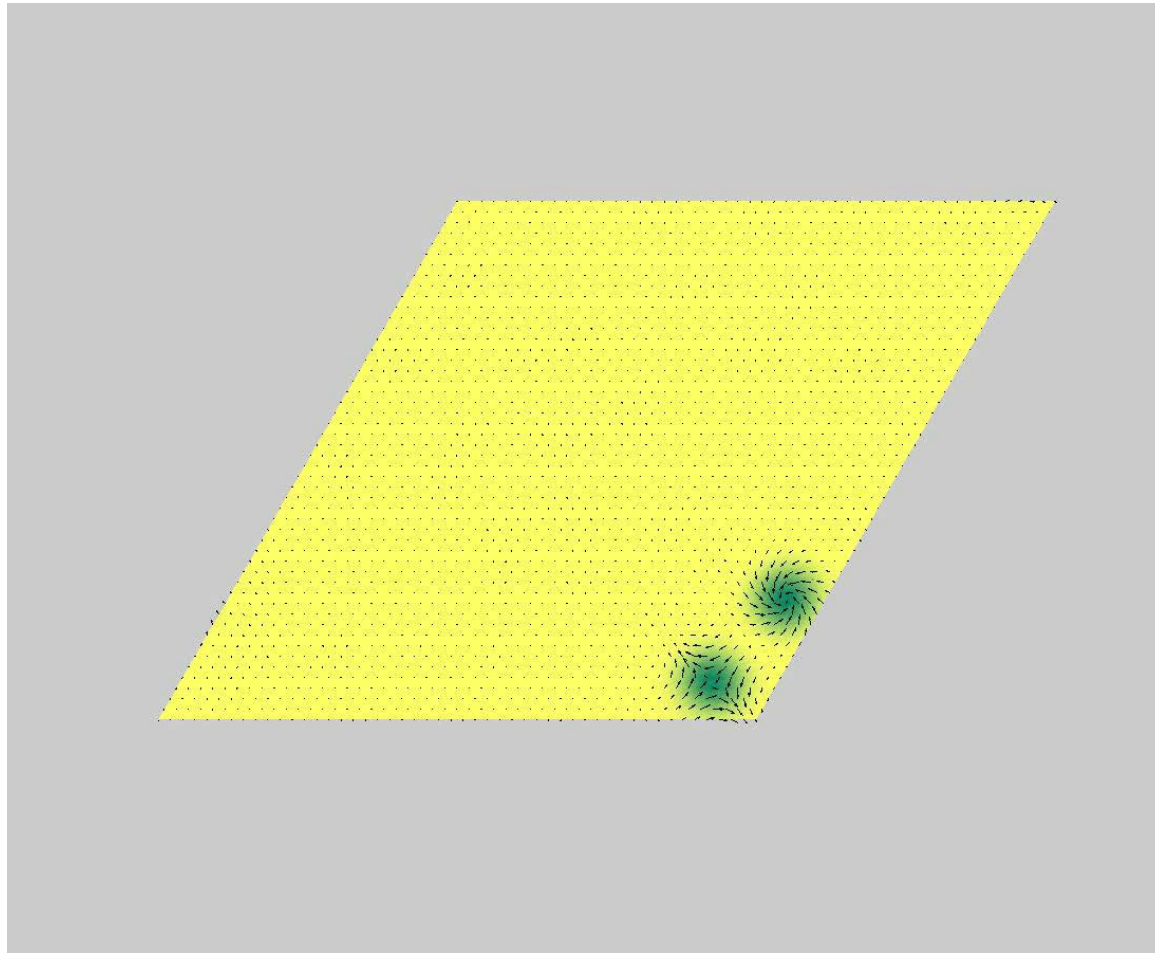
E_z -induced rotation of skyrmion-skyrmion pair



E_z -excitation of 3 skyrmions



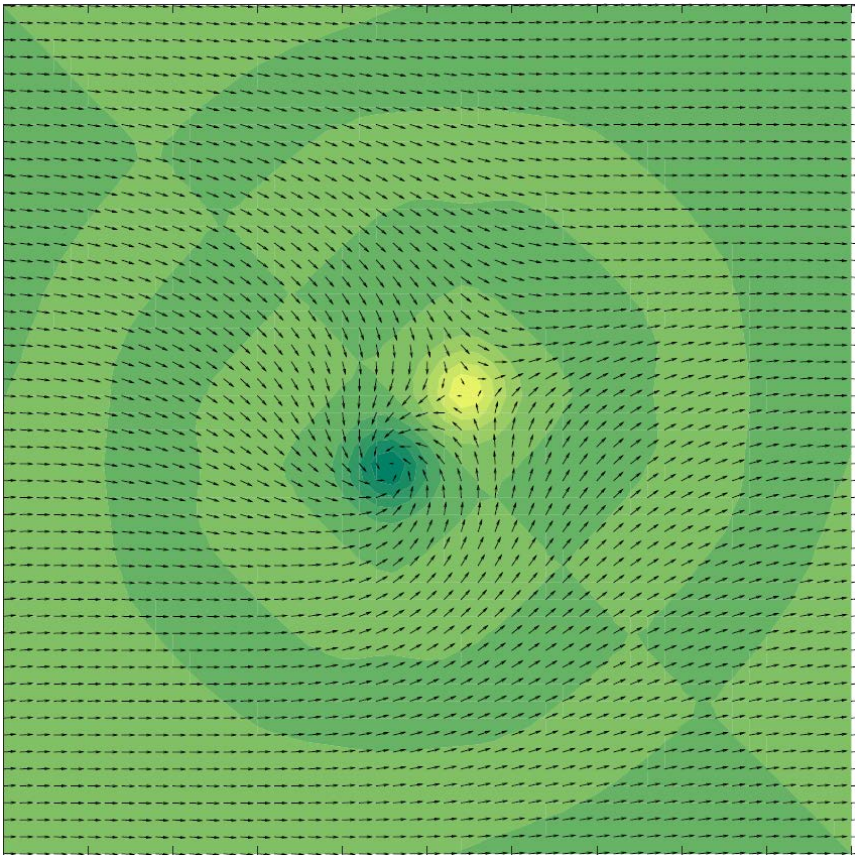
E_z -induced motion of skyrmion-antiskyrmion pair



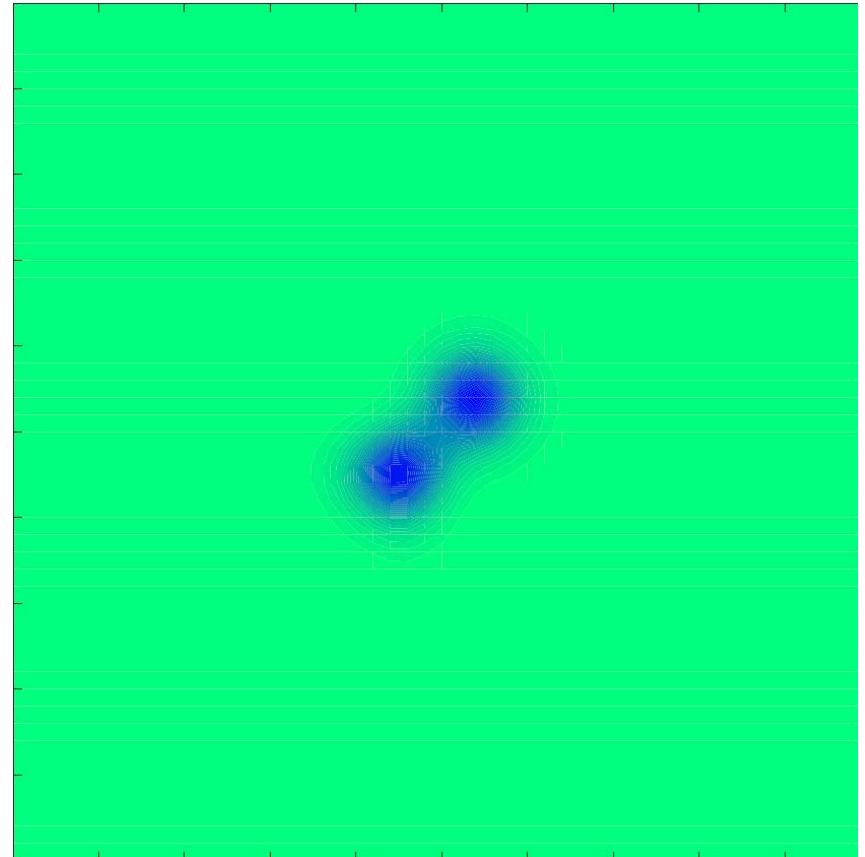
Bi-merons

easy plane anisotropy

Spin configuration



Topological density



Electric charge of magnetic vortex

MM PRL 96, 067601 (2006)

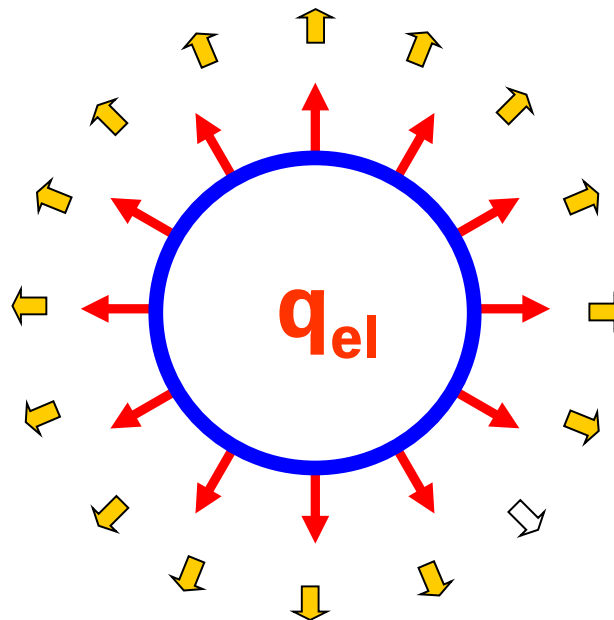
ME coupling

$$f_{me} = -g \mathbf{E} \cdot [\mathbf{m}(\nabla \cdot \mathbf{m}) - (\mathbf{m} \cdot \nabla)\mathbf{m}]$$

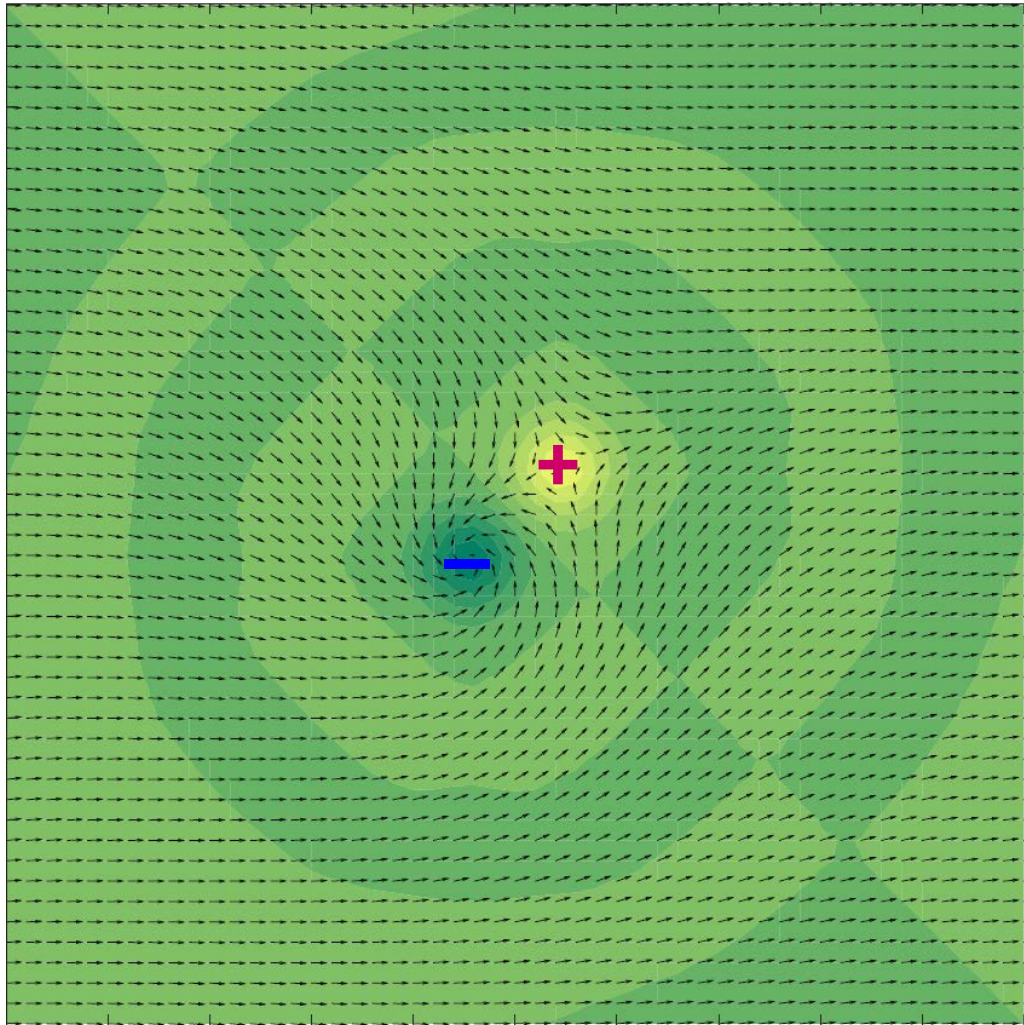
Force

$$\mathbf{F} = -\frac{\partial f_{me}}{\partial \mathbf{X}} = -g \mathbf{E} \oint_{\partial G} dx_i (\hat{\mathbf{z}} \cdot \mathbf{m} \times \partial_i \mathbf{m})$$

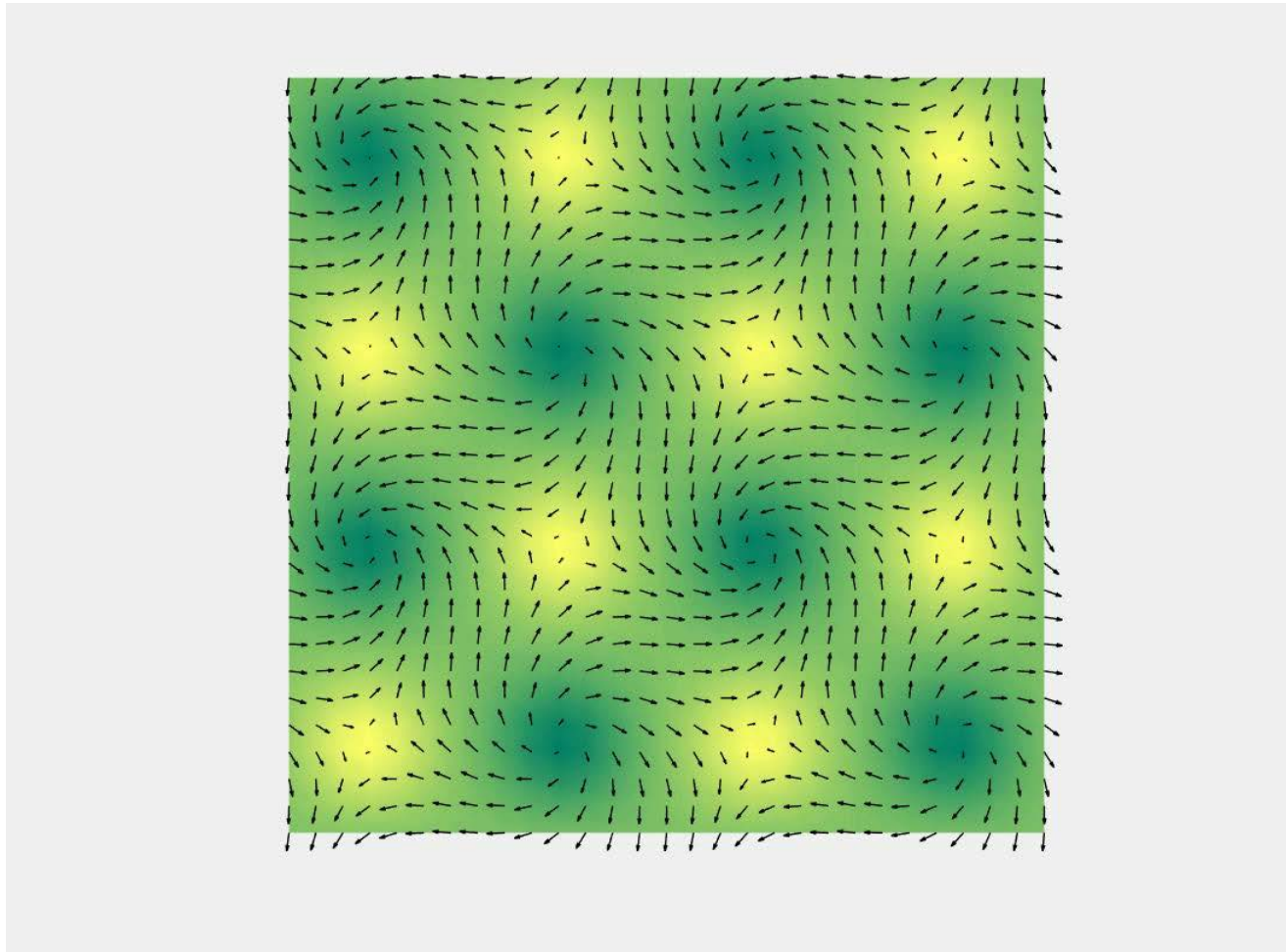
winding number



Electrically-charged merons



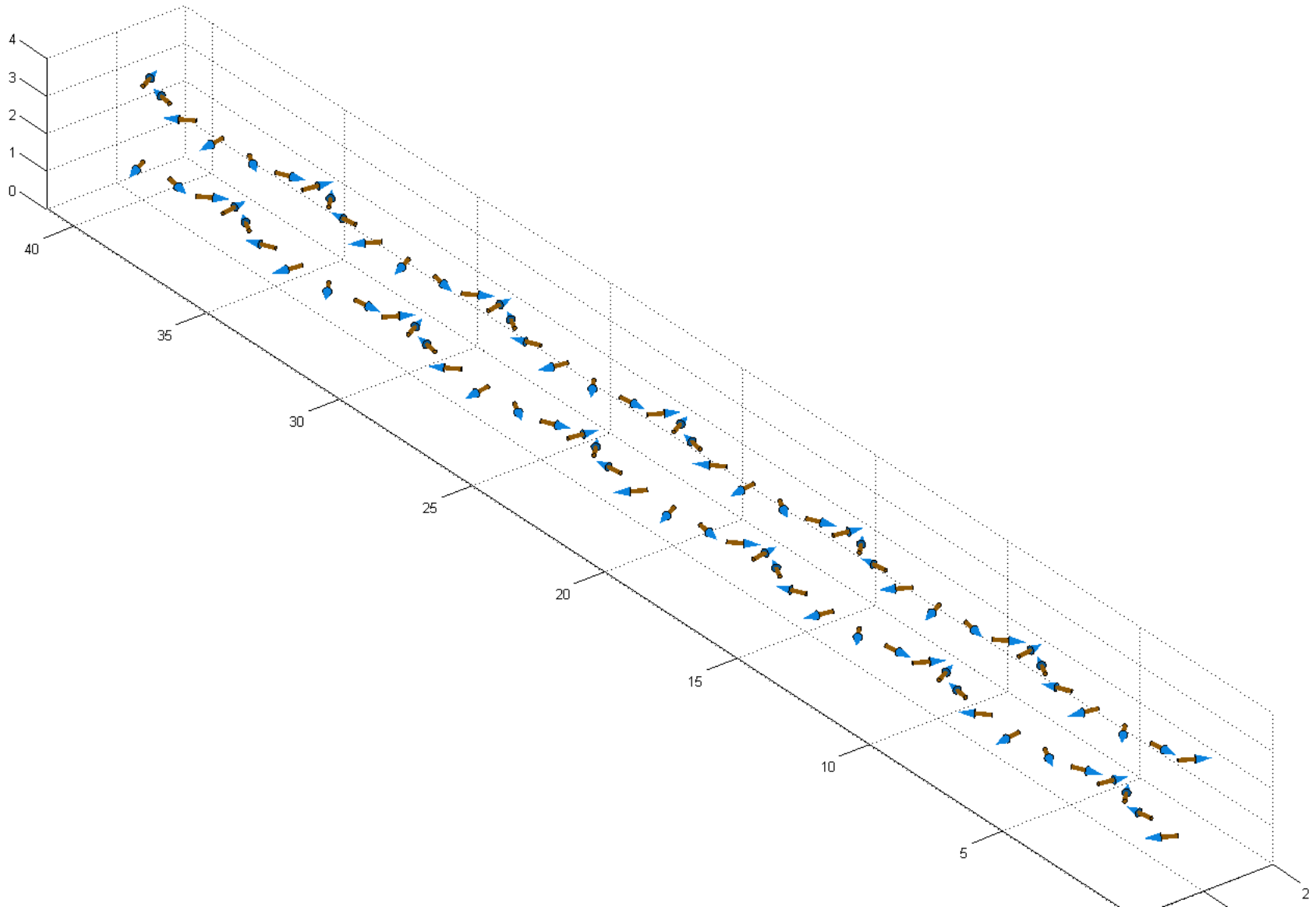
E_x excites plasma oscillations along y



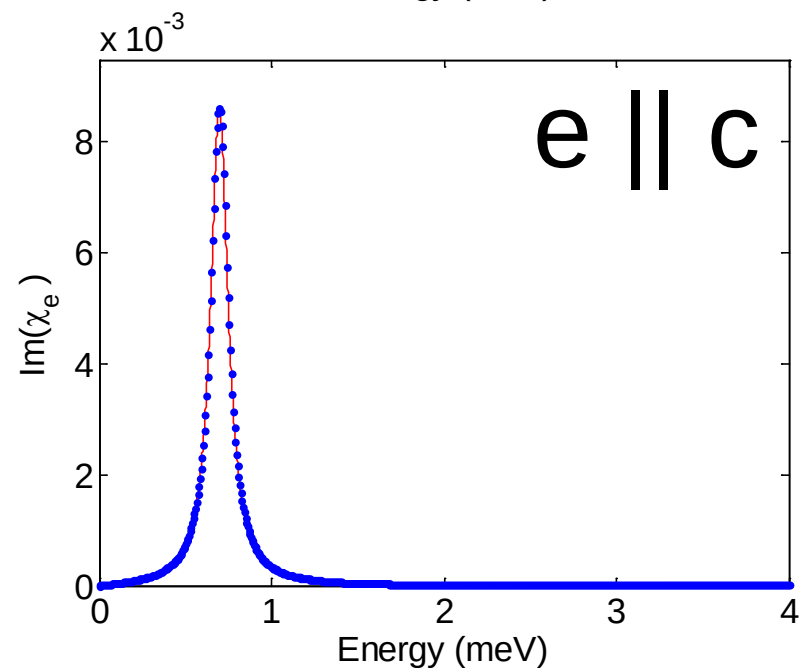
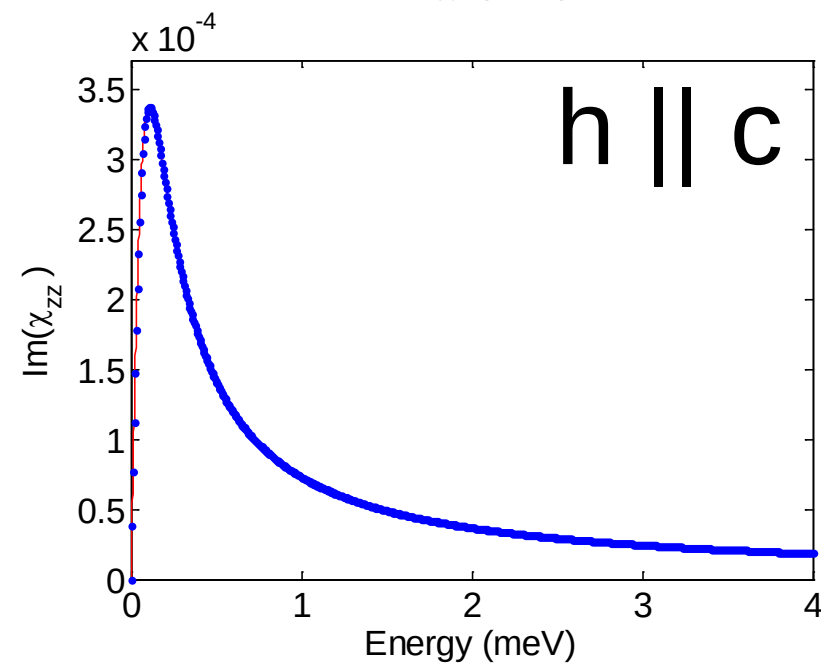
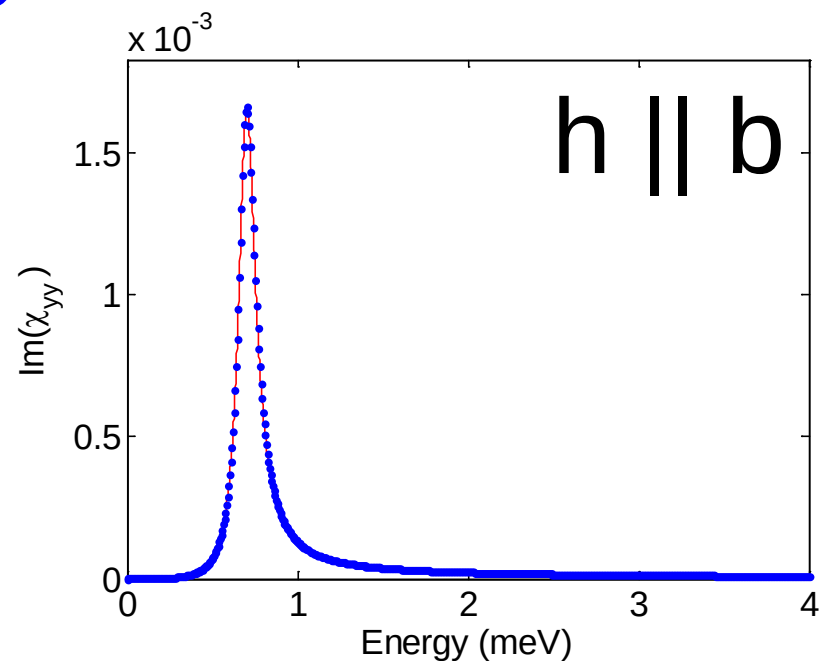
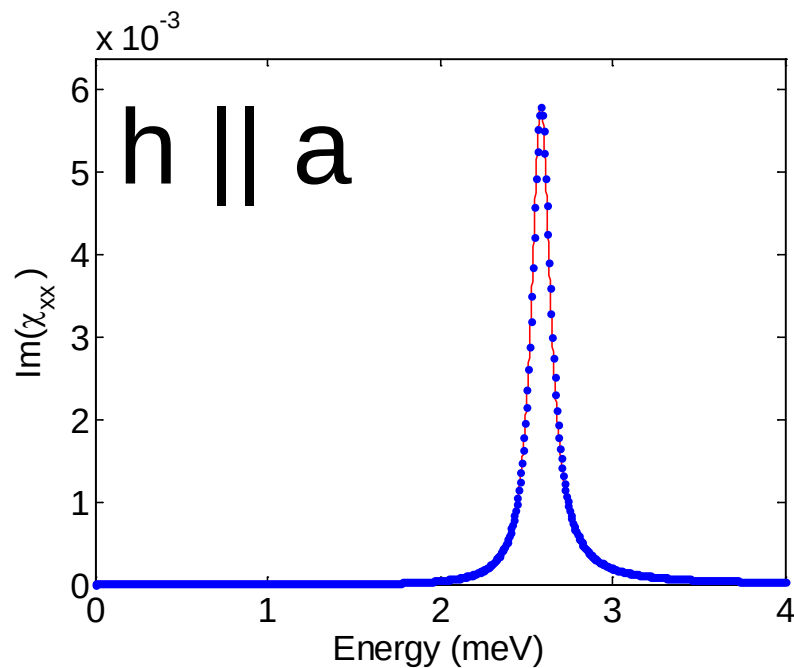
Summary

- Magnetoelectric coupling allows for electric field control of light propagation through non-collinear spin textures and excitation of dynamics of topological magnetic defects with an oscillating electric field.
- Low energy losses in Mott insulators.

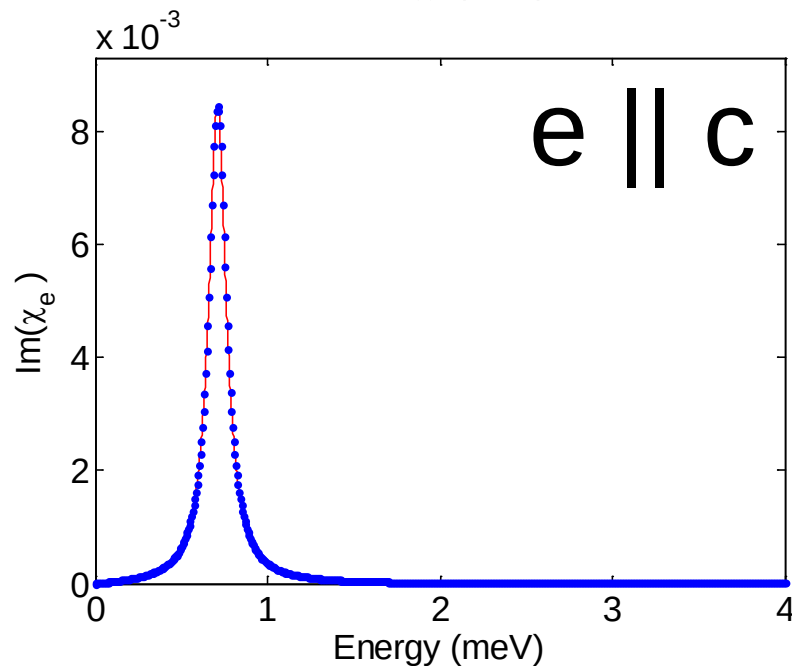
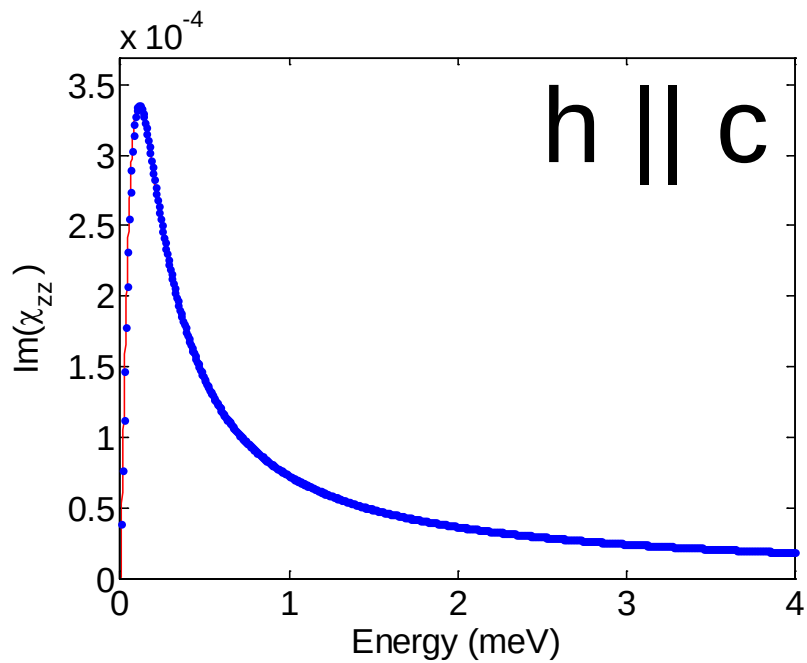
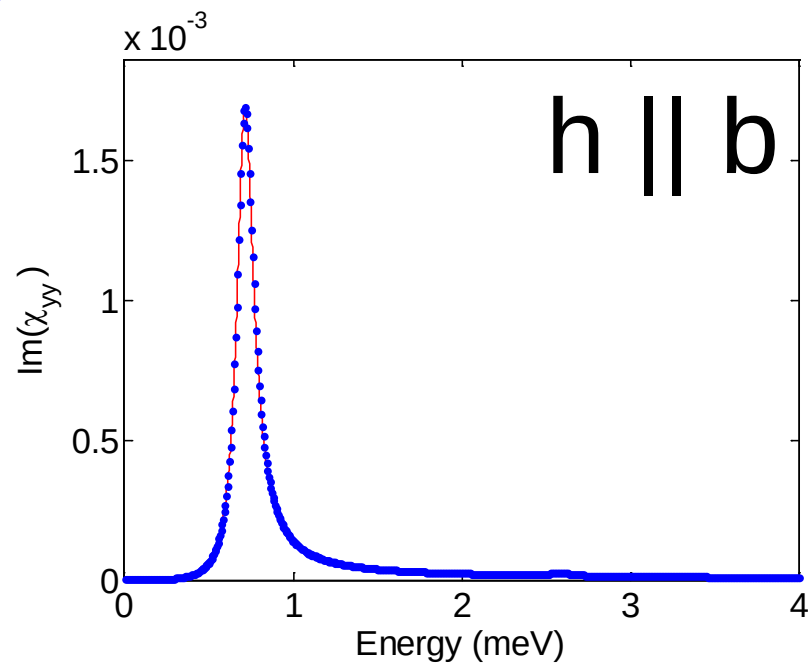
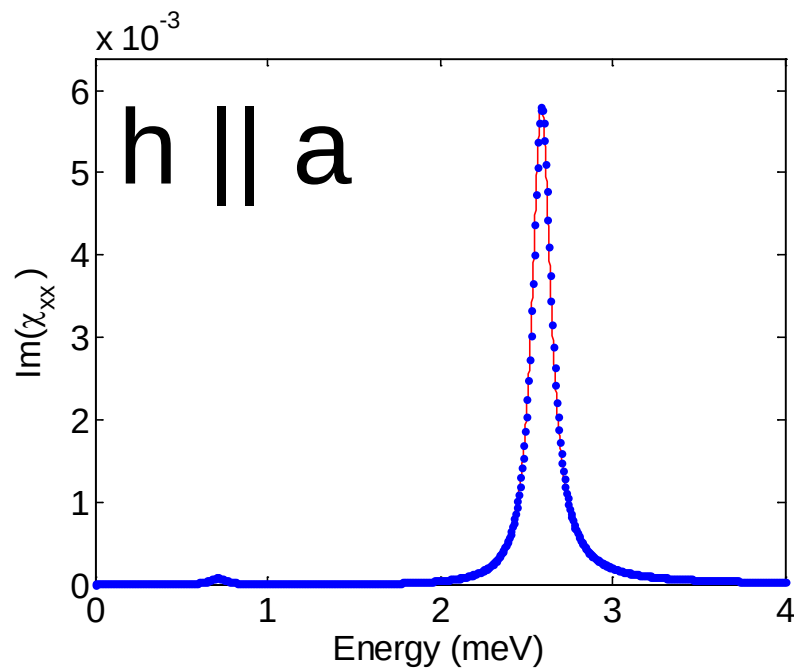
Spin configuration



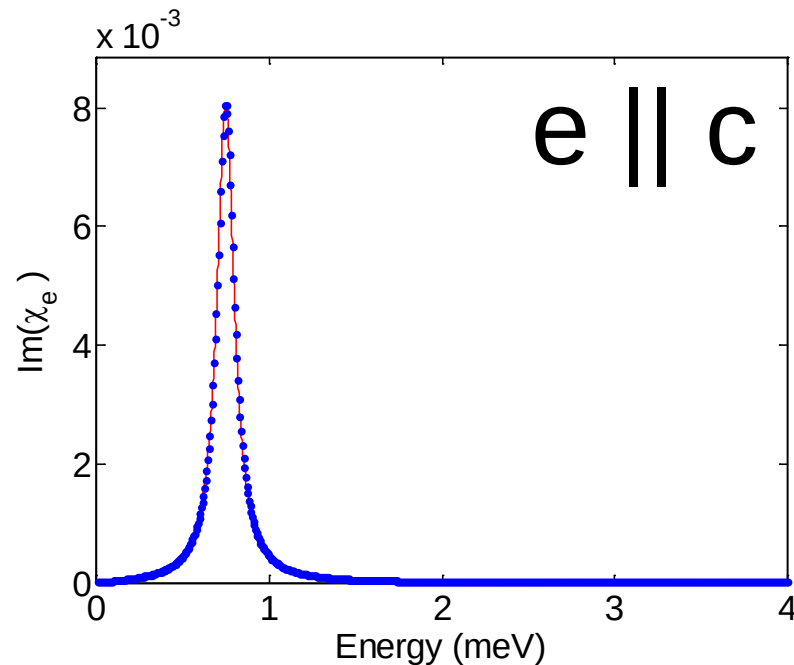
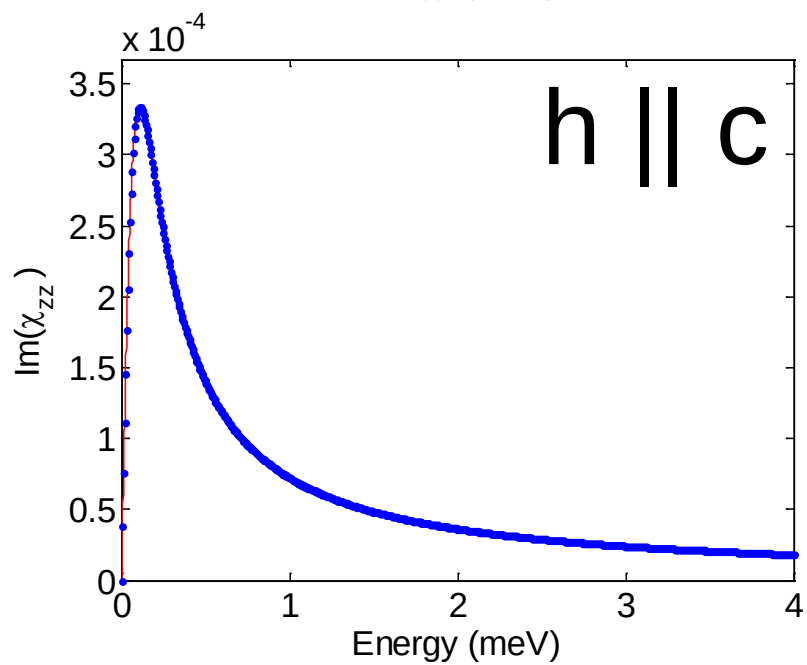
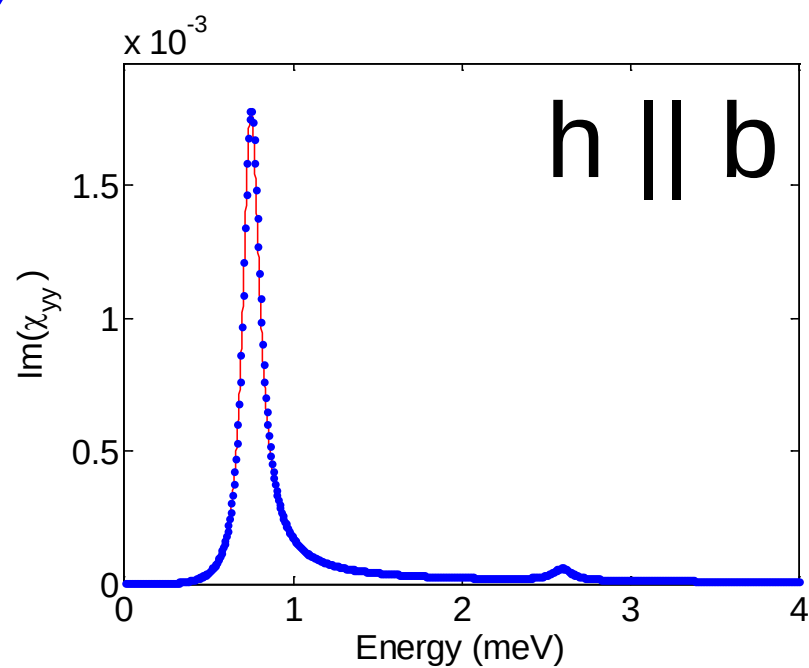
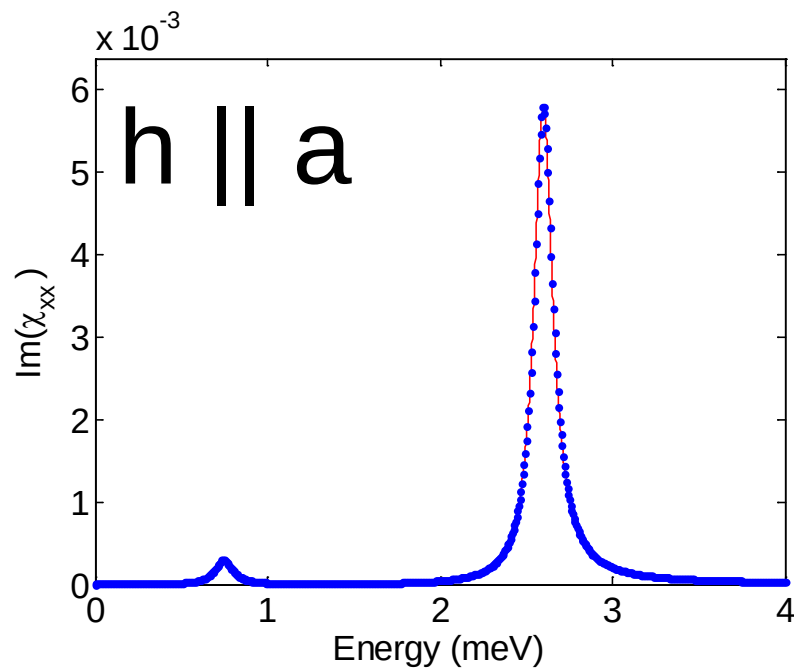
AFMR and E-magnons at $H = 0\text{T}$



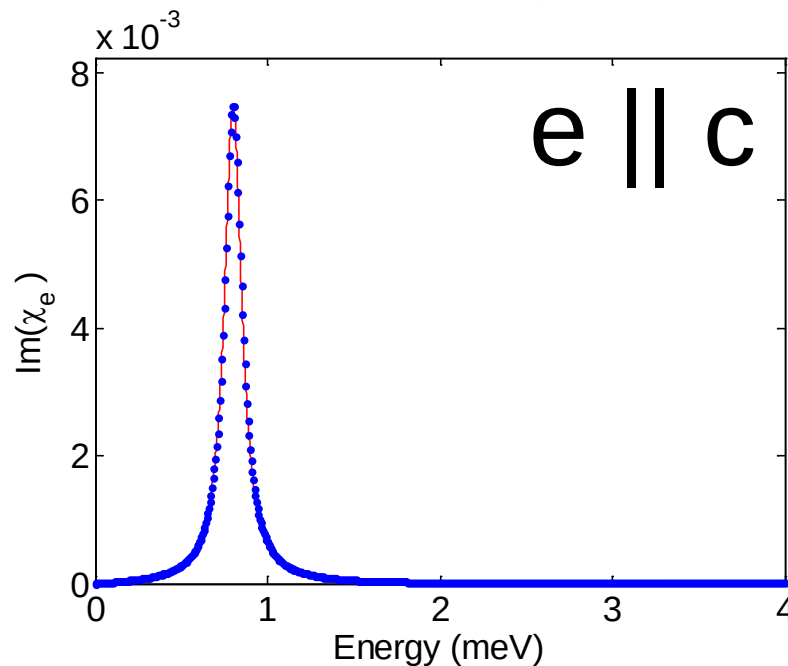
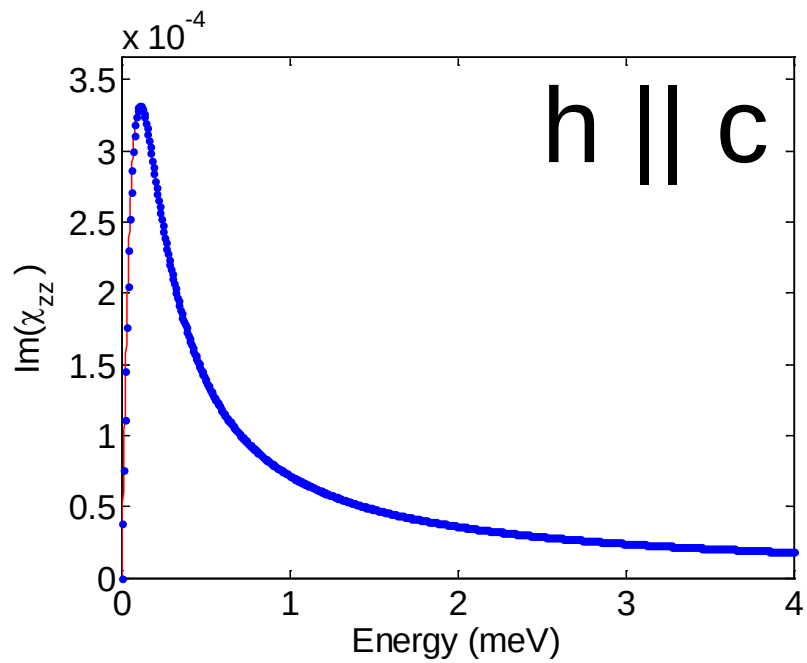
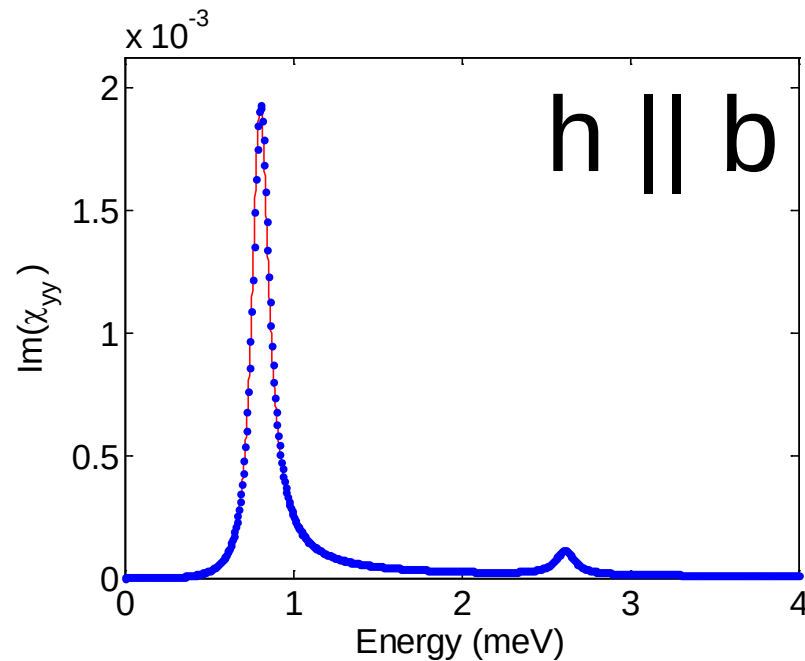
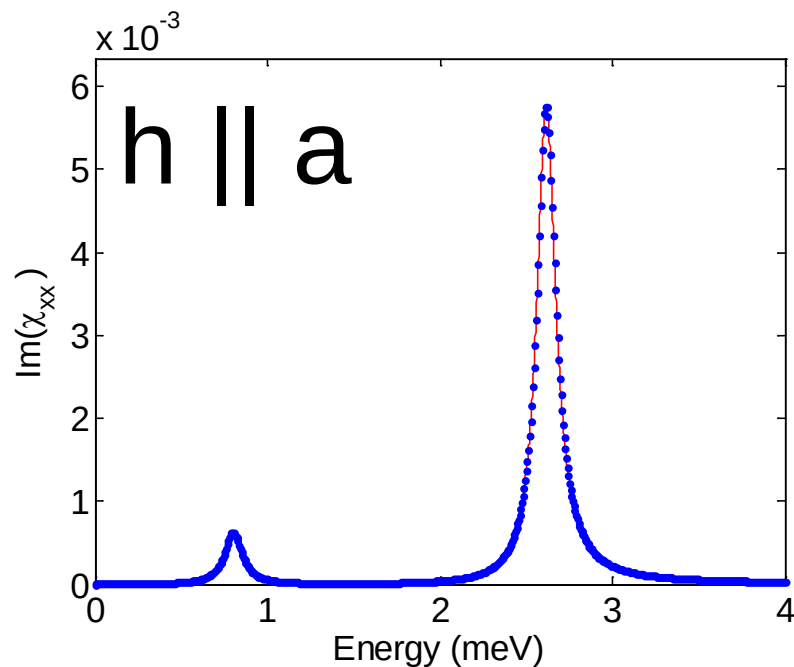
AFMR and E-magnons at $H = 1\text{T}$



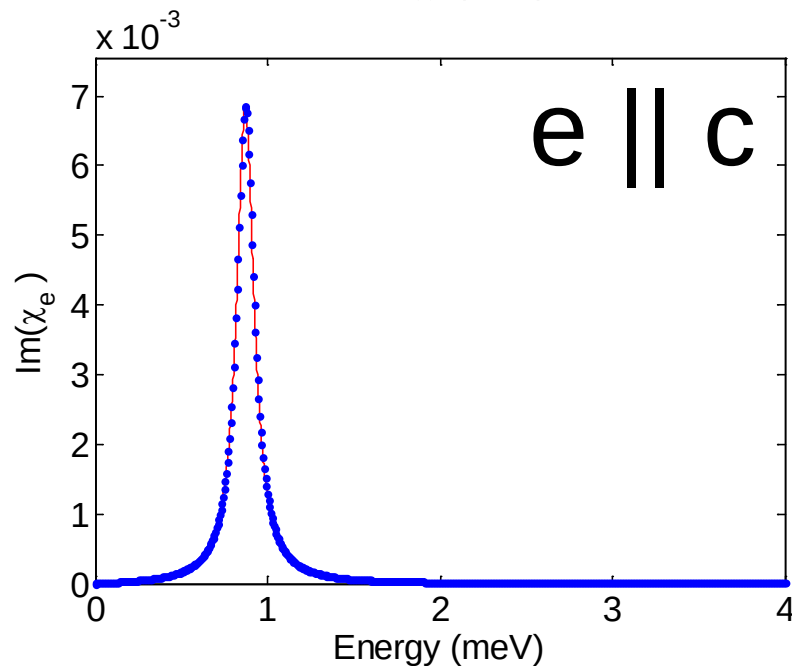
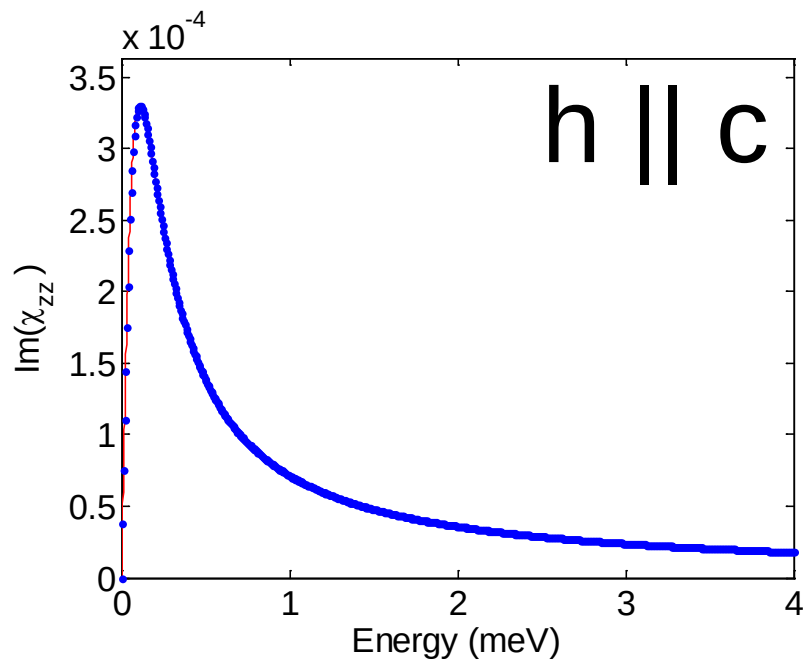
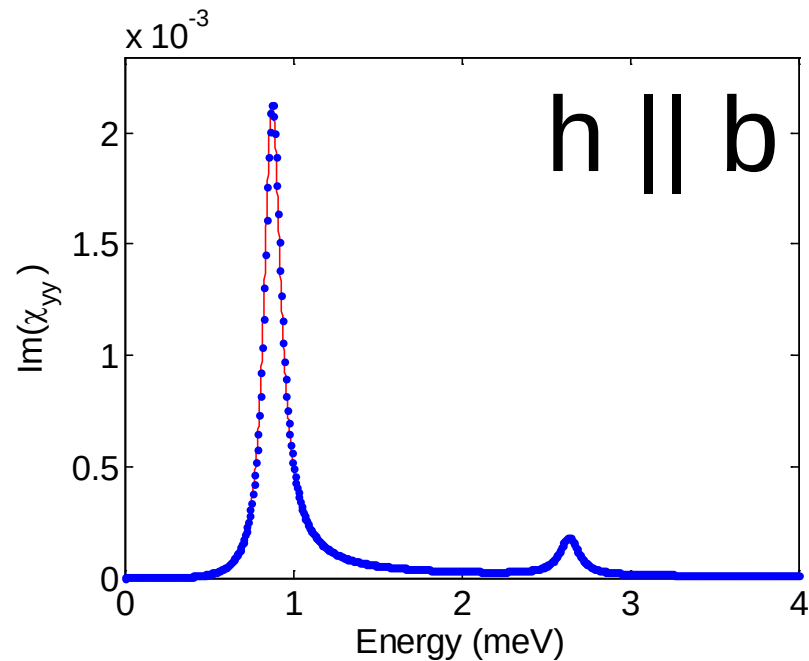
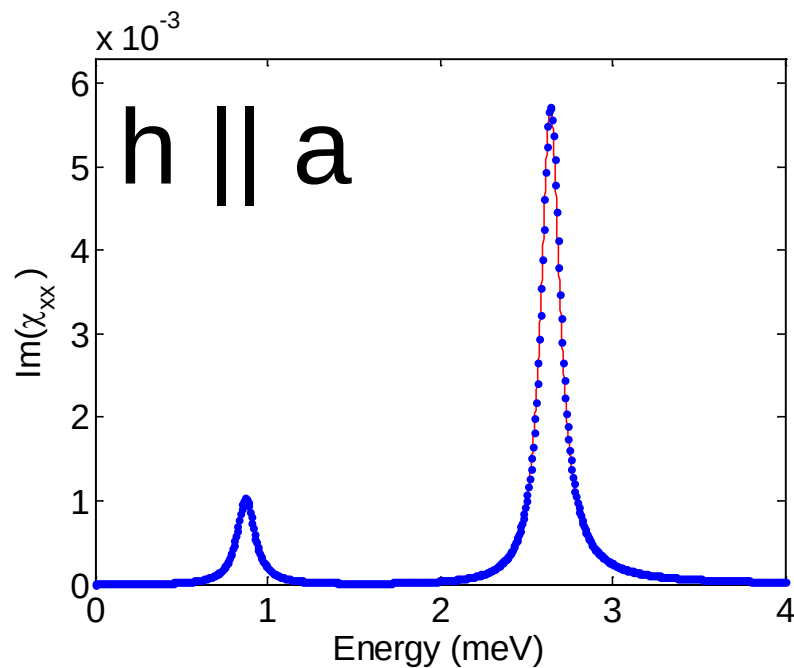
AFMR and E-magnons at $H = 2\text{T}$



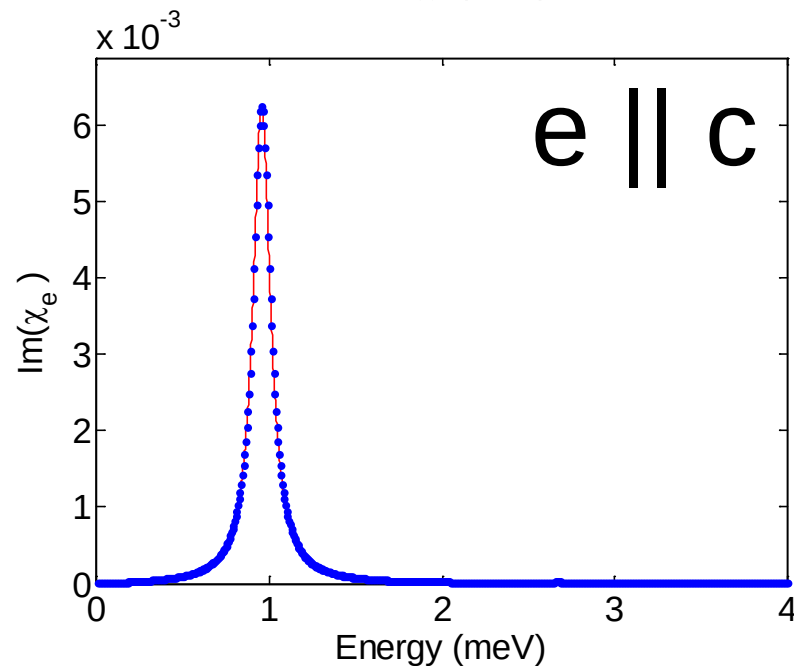
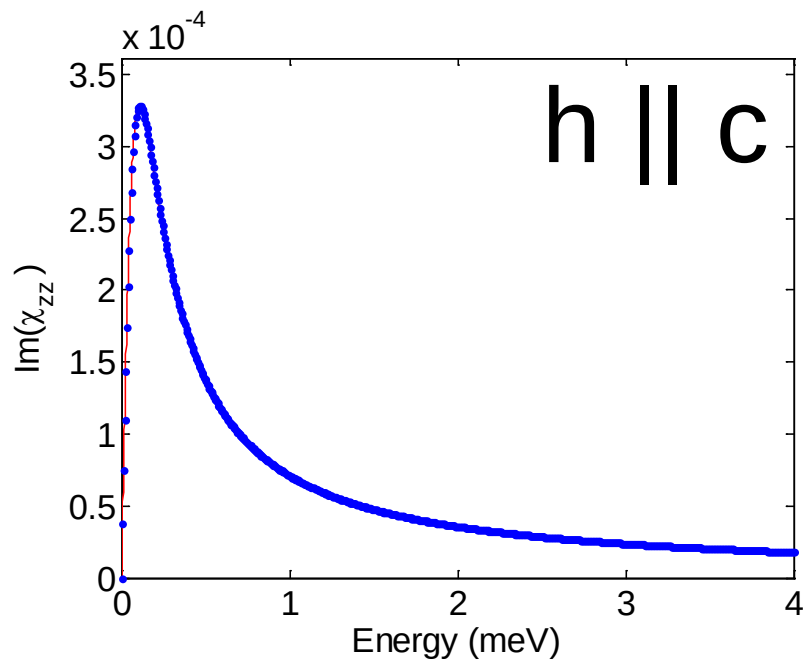
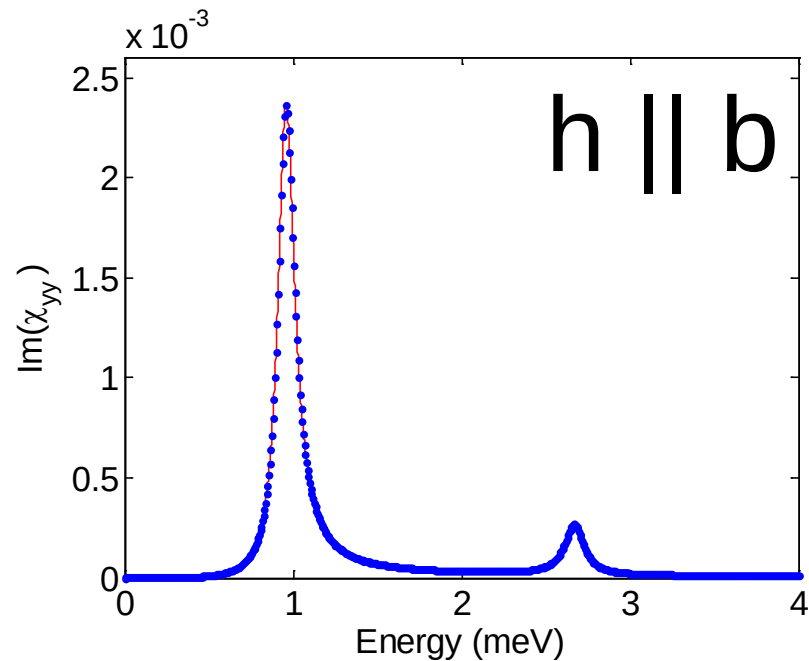
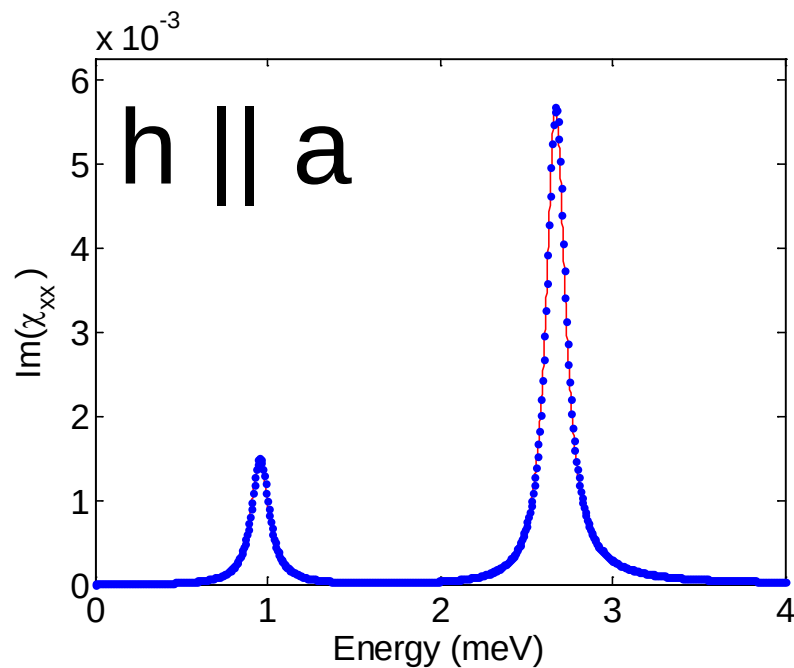
AFMR and E-magnons at $H = 3\text{T}$



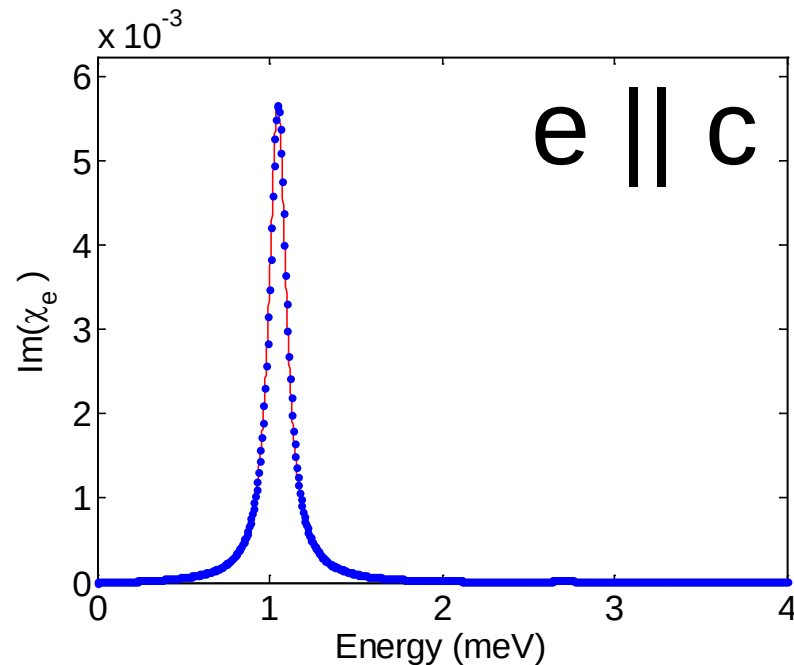
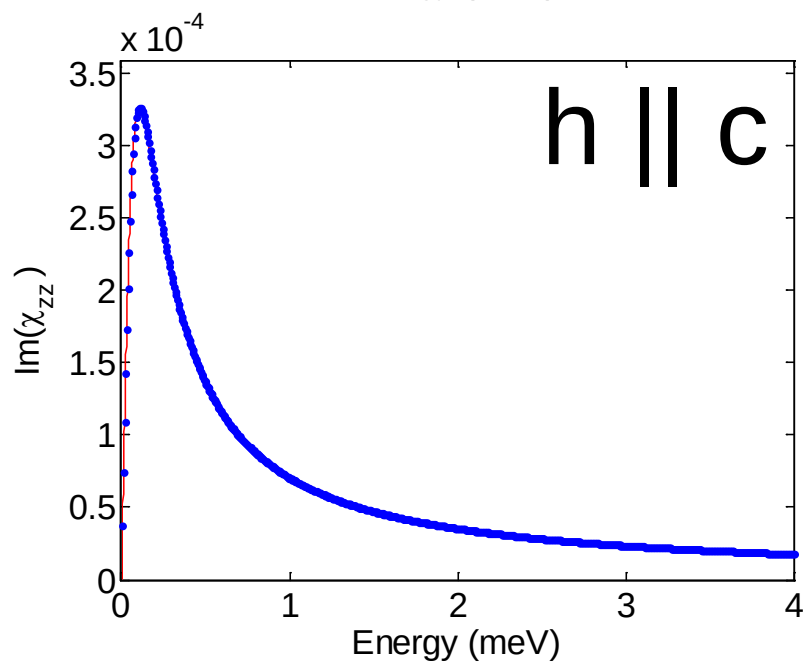
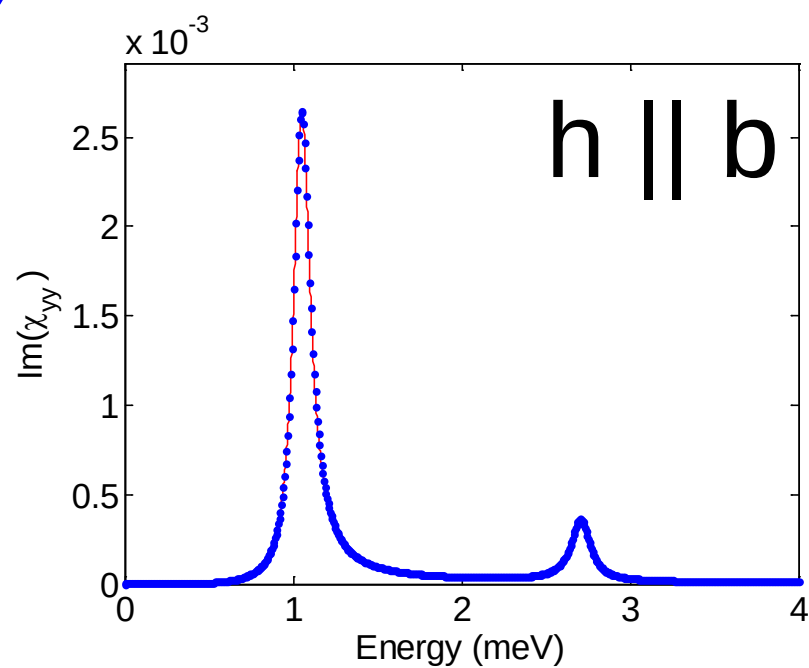
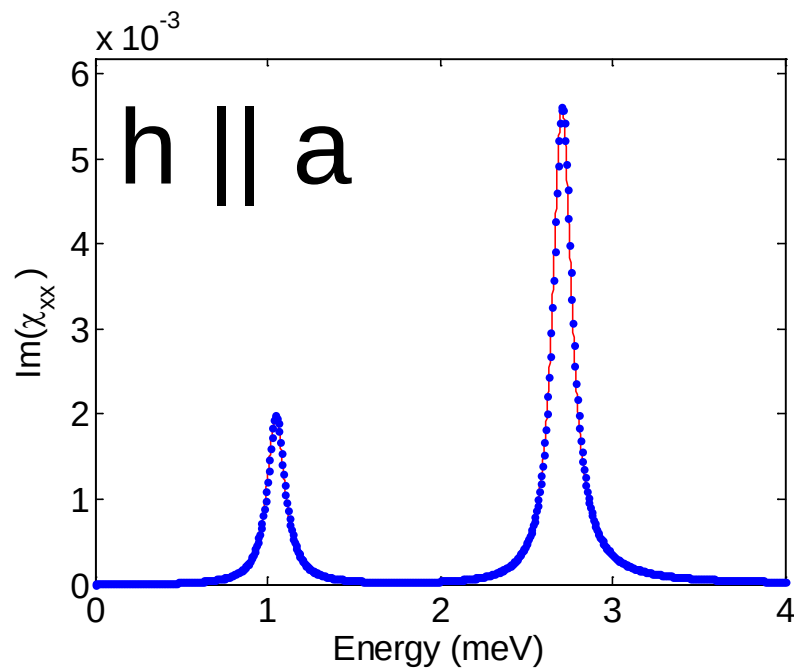
AFMR and E-magnons at $H = 4\text{T}$



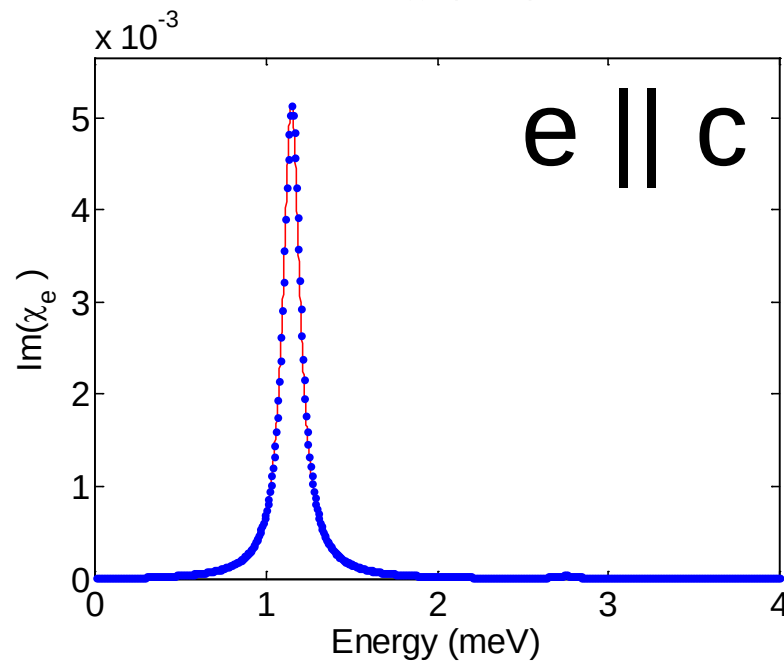
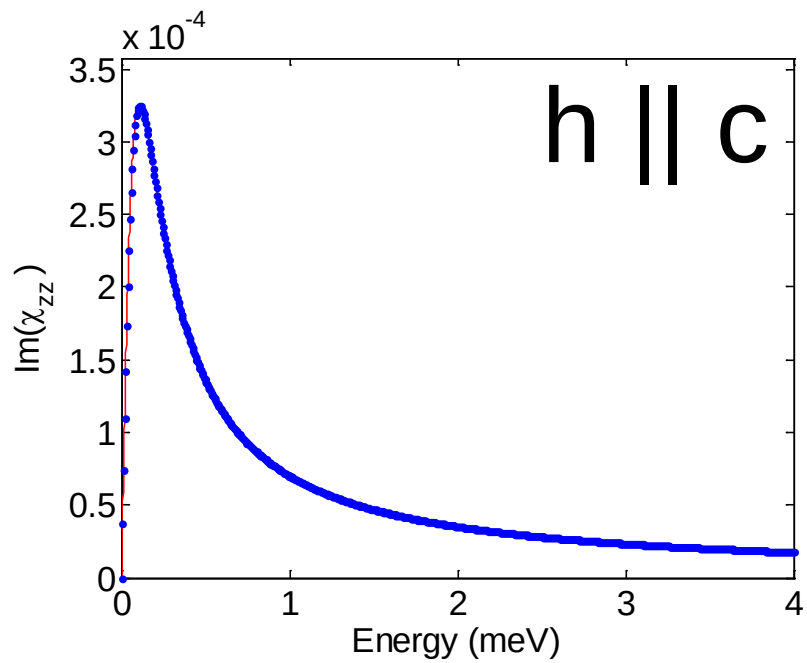
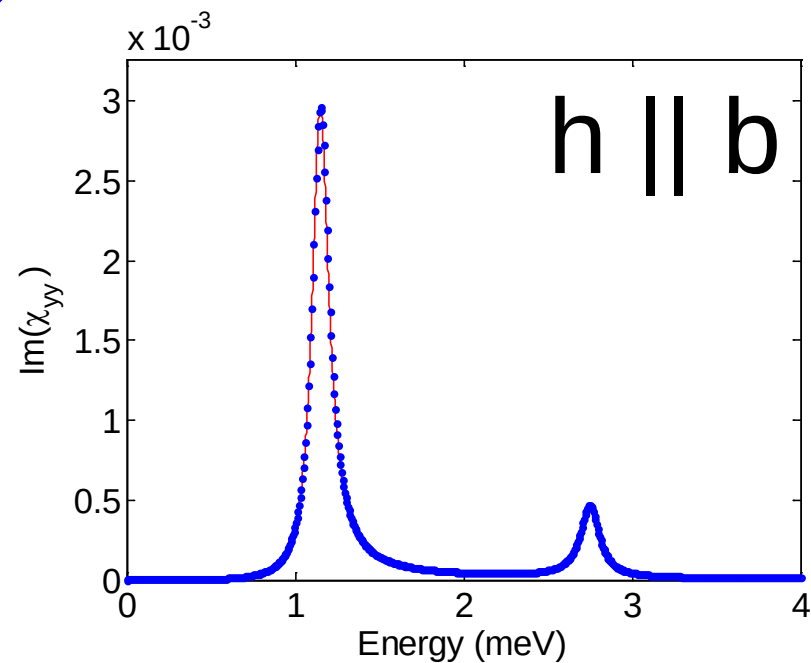
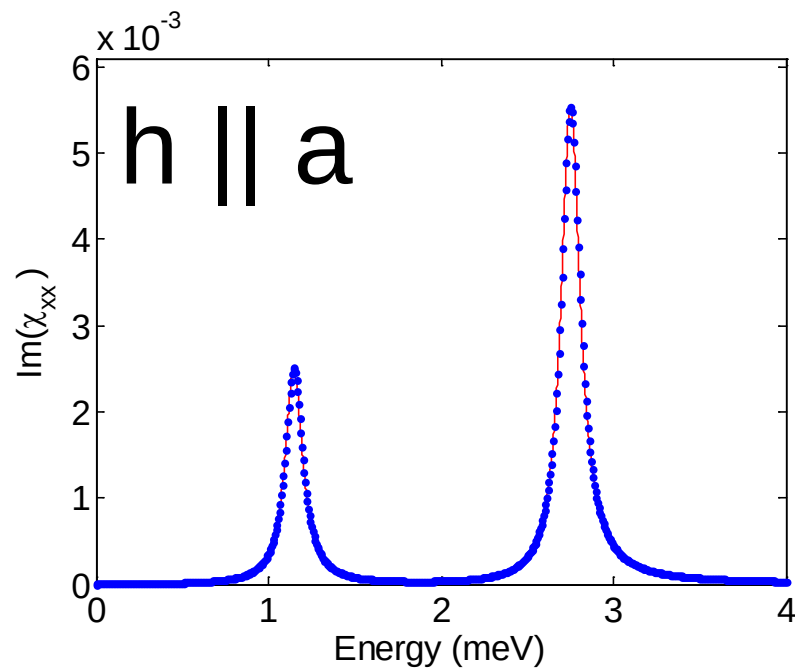
AFMR and E-magnons at $H = 5\text{T}$



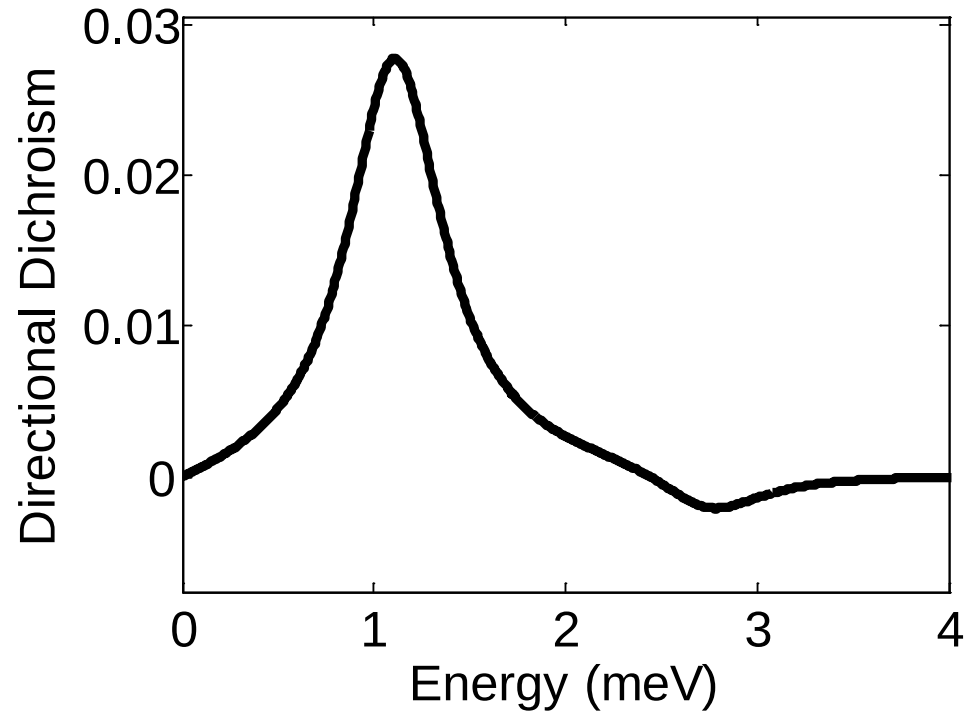
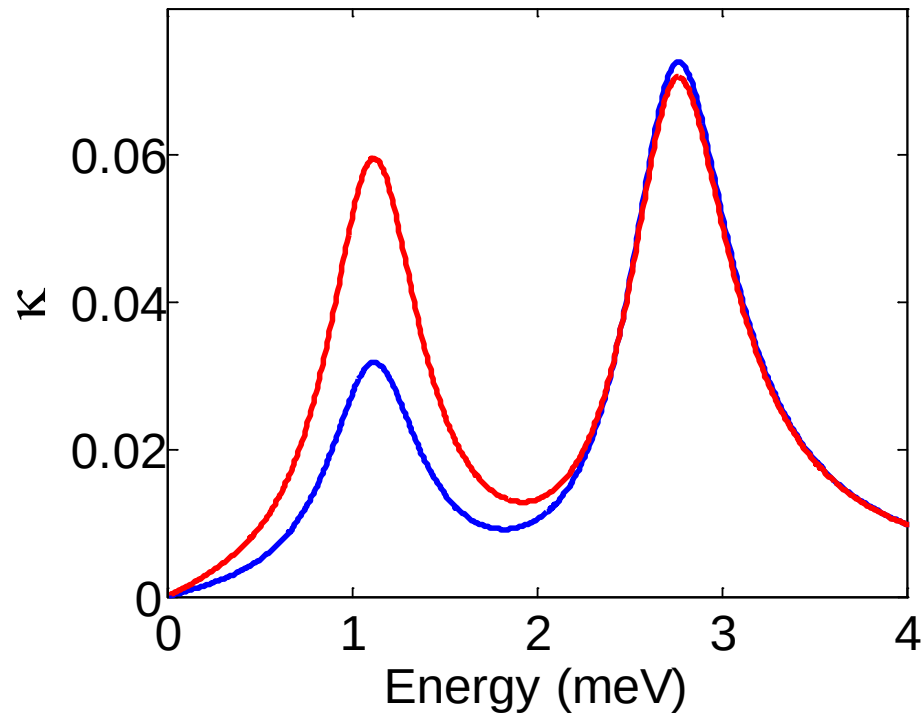
AFMR and E-magnons at $H = 6\text{T}$



AFMR and E-magnons at $H = 7\text{T}$



Extinction coefficient and directional dichroism at $H = 7\text{T}$



Spin flop

